

Is it worth the effort to find Lefschetz thimbles?

Dávid Pesznyák

in collaboration with
Attila Pásztor

based on Phys.Rev.D 113 (2026) 1, 014506



Eötvös Loránd
University

ELFT Particle Physics Seminar
February 24, 2026



HUN-REN Wigner
Research Centre for
Physics

Motivation: Complex Action Problem

partition function as a path integral

$$\mathcal{Z} = \int \mathcal{D}\phi e^{-S[\phi]} = \int \mathcal{D}\phi w[\phi]$$

Motivation: Complex Action Problem

partition function as a path integral

$$\mathcal{Z} = \int \mathcal{D}\phi e^{-S[\phi]} = \int \mathcal{D}\phi w[\phi]$$

if weights $w[\phi] \notin \mathbb{R}^+$, usual MCMC methods relying on importance sampling are not applicable:

Motivation: Complex Action Problem

partition function as a path integral

$$\mathcal{Z} = \int \mathcal{D}\phi e^{-S[\phi]} = \int \mathcal{D}\phi w[\phi]$$

if weights $w[\phi] \notin \mathbb{R}^+$, usual MCMC methods relying on importance sampling are not applicable:

Complex Action Problem

Motivation: Complex Action Problem

partition function as a path integral

$$\mathcal{Z} = \int \mathcal{D}\phi e^{-S[\phi]} = \int \mathcal{D}\phi w[\phi]$$

if weights $w[\phi] \notin \mathbb{R}^+$, usual MCMC methods relying on importance sampling are not applicable:

Complex Action Problem

complex action problem hinders studies in

Motivation: Complex Action Problem

partition function as a path integral

$$\mathcal{Z} = \int \mathcal{D}\phi e^{-S[\phi]} = \int \mathcal{D}\phi w[\phi]$$

if weights $w[\phi] \notin \mathbb{R}^+$, usual MCMC methods relying on importance sampling are not applicable:

Complex Action Problem

complex action problem hinders studies in

- finite-density QCD(-like models)

Motivation: Complex Action Problem

partition function as a path integral

$$\mathcal{Z} = \int \mathcal{D}\phi e^{-S[\phi]} = \int \mathcal{D}\phi w[\phi]$$

if weights $w[\phi] \notin \mathbb{R}^+$, usual MCMC methods relying on importance sampling are not applicable:

Complex Action Problem

complex action problem hinders studies in

- finite-density QCD(-like models)
- strongly correlated electron systems away from half-filling

Motivation: Complex Action Problem

partition function as a path integral

$$\mathcal{Z} = \int \mathcal{D}\phi e^{-S[\phi]} = \int \mathcal{D}\phi w[\phi]$$

if weights $w[\phi] \notin \mathbb{R}^+$, usual MCMC methods relying on importance sampling are not applicable:

Complex Action Problem

complex action problem hinders studies in

- finite-density QCD(-like models)
- strongly correlated electron systems away from half-filling
- Minkowskian time evolution in any quantum system

Motivation: Complex Action Problem

sign problem in reweighting

$$\langle \mathcal{O} \rangle_w = \frac{\int \mathcal{D}\phi \mathcal{O}[\phi] w[\phi]}{\int \mathcal{D}\phi w[\phi]} = \frac{\int \mathcal{D}\phi \mathcal{O}[\phi] \frac{w[\phi]}{r[\phi]} r[\phi]}{\int \mathcal{D}\phi \frac{w[\phi]}{r[\phi]} r[\phi]} = \frac{\langle \mathcal{O} \frac{w}{r} \rangle_r}{\langle \frac{w}{r} \rangle_r}$$

Motivation: Complex Action Problem

sign problem in reweighting

$$\langle \mathcal{O} \rangle_w = \frac{\int \mathcal{D}\phi \mathcal{O}[\phi] w[\phi]}{\int \mathcal{D}\phi w[\phi]} = \frac{\int \mathcal{D}\phi \mathcal{O}[\phi] \frac{w[\phi]}{r[\phi]} r[\phi]}{\int \mathcal{D}\phi \frac{w[\phi]}{r[\phi]} r[\phi]} = \frac{\langle \mathcal{O} \frac{w}{r} \rangle_r}{\langle \frac{w}{r} \rangle_r}$$

common choice is phase quenched: $r = |w| \implies w/r = e^{-i\text{Im}S} \equiv e^{i\theta}$

Motivation: Complex Action Problem

sign problem in reweighting

$$\langle \mathcal{O} \rangle_w = \frac{\int \mathcal{D}\phi \mathcal{O}[\phi] w[\phi]}{\int \mathcal{D}\phi w[\phi]} = \frac{\int \mathcal{D}\phi \mathcal{O}[\phi] \frac{w[\phi]}{r[\phi]} r[\phi]}{\int \mathcal{D}\phi \frac{w[\phi]}{r[\phi]} r[\phi]} = \frac{\langle \mathcal{O} \frac{w}{r} \rangle_r}{\langle \frac{w}{r} \rangle_r}$$

common choice is phase quenched: $r = |w| \implies w/r = e^{-i\text{Im}S} \equiv e^{i\theta}$

- fluctuations in $e^{i\theta} \implies$ large cancellations \implies large uncertainties

Motivation: Complex Action Problem

sign problem in reweighting

$$\langle \mathcal{O} \rangle_w = \frac{\int \mathcal{D}\phi \mathcal{O}[\phi] w[\phi]}{\int \mathcal{D}\phi w[\phi]} = \frac{\int \mathcal{D}\phi \mathcal{O}[\phi] \frac{w[\phi]}{r[\phi]} r[\phi]}{\int \mathcal{D}\phi \frac{w[\phi]}{r[\phi]} r[\phi]} = \frac{\langle \mathcal{O} \frac{w}{r} \rangle_r}{\langle \frac{w}{r} \rangle_r}$$

common choice is phase quenched: $r = |w| \implies w/r = e^{-i\text{Im}S} \equiv e^{i\theta}$

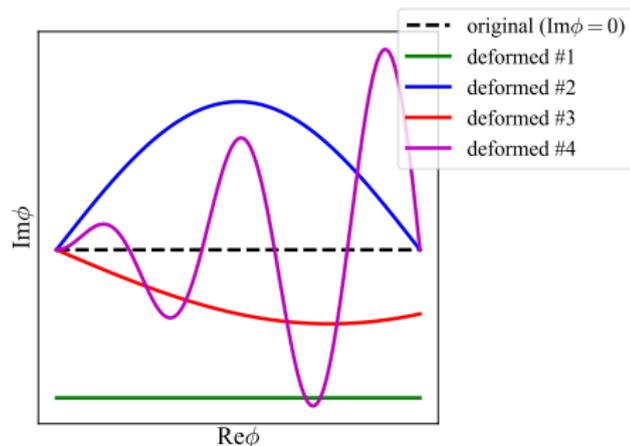
- fluctuations in $e^{i\theta} \implies$ large cancellations \implies large uncertainties
- severity of the sign problem:

$$\left\langle \frac{w}{r} \right\rangle_r = \frac{\mathcal{Z}_w}{\mathcal{Z}_r} = \sigma \implies \begin{cases} 1 & \sim \text{perfect!} \\ \approx 0 & \sim \text{not so much...} \end{cases}$$

Fighting the sign problem?

N -dimensional integral over real fields to

- $2N$ -dimensional stochastic dynamics in the real and imaginary parts of complexified fields (i.e. *Complex Langevin Dynamics*)
- N -dimensional integral with deformed integration contour/manifold into the complexified field space



Contour deformations

assuming that \mathcal{M} and \mathcal{M}' are in the same homology class

$$\mathcal{Z} = \int_{\mathcal{M}} \mathcal{D}\phi w[\phi] = \int_{\mathcal{M}'} \mathcal{D}\phi' w[\phi']$$

Contour deformations

assuming that \mathcal{M} and \mathcal{M}' are in the same homology class

$$\mathcal{Z} = \int_{\mathcal{M}} \mathcal{D}\phi w[\phi] = \int_{\mathcal{M}'} \mathcal{D}\phi' w[\phi']$$

parameterize \mathcal{M}' by the original real variable ϕ

$$\mathcal{Z} = \int_{\mathcal{M}} \mathcal{D}\phi \det \mathcal{J}(\phi) w[\phi'(\phi)] = \int_{\mathcal{M}} \mathcal{D}\phi e^{-S^{\text{eff}}[\phi]}$$

Contour deformations

assuming that \mathcal{M} and \mathcal{M}' are in the same homology class

$$\mathcal{Z} = \int_{\mathcal{M}} \mathcal{D}\phi w[\phi] = \int_{\mathcal{M}'} \mathcal{D}\phi' w[\phi']$$

parameterize \mathcal{M}' by the original real variable ϕ

$$\mathcal{Z} = \int_{\mathcal{M}} \mathcal{D}\phi \det \mathcal{J}(\phi) w[\phi'(\phi)] = \int_{\mathcal{M}} \mathcal{D}\phi e^{-S^{\text{eff}}[\phi]}$$

phase-quenched partition function

$$\mathcal{Z}_{\text{PQ}}^{\text{def}} = \int_{\mathcal{M}} \mathcal{D}\phi \left| \det \mathcal{J}(\phi) w[\phi'(\phi)] \right| = \int_{\mathcal{M}} \mathcal{D}\phi e^{-S_{\text{PQ}}^{\text{eff}}[\phi]}$$

Contour deformations: improving the sign problem

- \mathcal{Z} is invariant (physics unchanged)

Contour deformations: improving the sign problem

- \mathcal{Z} is invariant (physics unchanged)
- but $\mathcal{Z}_{PQ}^{\text{def}}$ does change!

Contour deformations: improving the sign problem

- \mathcal{Z} is invariant (physics unchanged)
- but $\mathcal{Z}_{\text{PQ}}^{\text{def}}$ does change!

$$\sigma = \langle e^{i\theta} \rangle_{\text{PQ}}^{\text{def}} = \frac{\mathcal{Z}}{\mathcal{Z}_{\text{PQ}}^{\text{def}}} = \left\langle \frac{\det \mathcal{J} w}{|\det \mathcal{J} w|} \right\rangle_{\text{PQ}}^{\text{def}}$$

Contour deformations: improving the sign problem

- \mathcal{Z} is invariant (physics unchanged)
- but $\mathcal{Z}_{\text{PQ}}^{\text{def}}$ does change!

$$\sigma = \langle e^{i\theta} \rangle_{\text{PQ}}^{\text{def}} = \frac{\mathcal{Z}}{\mathcal{Z}_{\text{PQ}}^{\text{def}}} = \left\langle \frac{\det \mathcal{J} w}{|\det \mathcal{J} w|} \right\rangle_{\text{PQ}}^{\text{def}}$$

finding a good manifold \mathcal{M}' leads to a milder sign problem

Three methods compared

- **Lefschetz thimbles:** generalized steepest-descent manifolds attached to critical points of the action

Three methods compared

- **Lefschetz thimbles:** generalized steepest-descent manifolds attached to critical points of the action
- **holomorphic flow:** continuous deformation towards thimbles via the holomorphic flow equations

Three methods compared

- **Lefschetz thimbles:** generalized steepest-descent manifolds attached to critical points of the action
- **holomorphic flow:** continuous deformation towards thimbles via the holomorphic flow equations
- **numerical sign optimization:** choose a tractable ansatz for $\phi'(\phi, \{p\})$ and optimize finite set of parameters $\{p\}$

Three methods compared

- **Lefschetz thimbles:** generalized steepest-descent manifolds attached to critical points of the action
- **holomorphic flow:** continuous deformation towards thimbles via the holomorphic flow equations
- **numerical sign optimization:** choose a tractable ansatz for $\phi'(\phi, \{p\})$ and optimize finite set of parameters $\{p\}$

Is the thimble limit numerically optimal?

Method I: Lefschetz thimbles

critical points ϕ_c solve

$$\left. \frac{\partial S}{\partial \phi'} \right|_{\phi_c} = 0$$

Method I: Lefschetz thimbles

critical points ϕ_c solve

$$\left. \frac{\partial S}{\partial \phi'} \right|_{\phi_c} = 0$$

thimbles are defined as a set of initial conditions of the downward flow in complexified space

$$\frac{d\phi'}{dt} = - \left(\frac{\partial S}{\partial \phi'} \right)^*$$

such that $\phi'(t \rightarrow \infty) = \phi_c$

Method I: Lefschetz thimbles

critical points ϕ_c solve

$$\left. \frac{\partial S}{\partial \phi'} \right|_{\phi_c} = 0$$

thimbles are defined as a set of initial conditions of the downward flow in complexified space

$$\frac{d\phi'}{dt} = - \left(\frac{\partial S}{\partial \phi'} \right)^*$$

such that $\phi'(t \rightarrow \infty) = \phi_c$

a disjoint thimble can be attached to each critical point of the action

Method I: Lefschetz thimbles

critical points ϕ_c solve

$$\left. \frac{\partial S}{\partial \phi'} \right|_{\phi_c} = 0$$

thimbles are defined as a set of initial conditions of the downward flow in complexified space

$$\frac{d\phi'}{dt} = - \left(\frac{\partial S}{\partial \phi'} \right)^*$$

such that $\phi'(t \rightarrow \infty) = \phi_c$

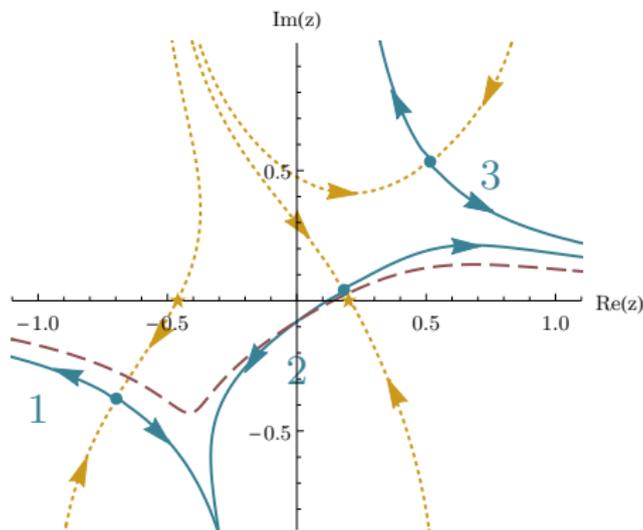
a disjoint thimble can be attached to each critical point of the action

on a given thimble, $\text{Im}S$ is constant \implies no phase fluctuations from e^{-S}

Method I: Lefschetz thimble (example)

- example with action from [arXiv:hep-lat/2007.05436]

$$S(z) = \frac{1}{G}z^2 - \log[(p^2 + i\mu)^2 + (z + m)^2]$$



Method I: Lefschetz thimbles

with thimble decomposition there are two remaining sources of cancellations

Method I: Lefschetz thimbles

with thimble decomposition there are two remaining sources of cancellations

- **residual phase** from the Jacobian

$$\mathcal{Z} = \int_{\mathcal{M}} \mathcal{D}\phi \det \mathcal{J}(\phi) e^{-S[\phi]}, \quad \det \mathcal{J}(\phi) \in \mathbb{C}$$

Method I: Lefschetz thimbles

with thimble decomposition there are two remaining sources of cancellations

- **residual phase** from the Jacobian

$$\mathcal{Z} = \int_{\mathcal{M}} \mathcal{D}\phi \det \mathcal{J}(\phi) e^{-S[\phi]}, \quad \det \mathcal{J}(\phi) \in \mathbb{C}$$

- **global sign problem** when more than one thimble contributes,

$$\mathcal{Z} = \sum_a n_a e^{-i\text{Im}S[\phi_c^a]} \mathcal{Z}_a, \quad \mathcal{Z}_a = \int_{\mathcal{T}_a} \mathcal{D}\phi' e^{-\text{Re}S[\phi']},$$

and different constant phases can lead to cancellations

Method I: Lefschetz thimbles

with thimble decomposition there are two remaining sources of cancellations

- **residual phase** from the Jacobian

$$\mathcal{Z} = \int_{\mathcal{M}} \mathcal{D}\phi \det \mathcal{J}(\phi) e^{-S[\phi]}, \quad \det \mathcal{J}(\phi) \in \mathbb{C}$$

- **global sign problem** when more than one thimble contributes,

$$\mathcal{Z} = \sum_a n_a e^{-i\text{Im}S[\phi_c^a]} \mathcal{Z}_a, \quad \mathcal{Z}_a = \int_{\mathcal{T}_a} \mathcal{D}\phi' e^{-\text{Re}S[\phi']},$$

and different constant phases can lead to cancellations

1. thimbles are clean analytically, but not necessarily optimal numerically

Method I: Lefschetz thimbles

with thimble decomposition there are two remaining sources of cancellations

- **residual phase** from the Jacobian

$$\mathcal{Z} = \int_{\mathcal{M}} \mathcal{D}\phi \det \mathcal{J}(\phi) e^{-S[\phi]}, \quad \det \mathcal{J}(\phi) \in \mathbb{C}$$

- **global sign problem** when more than one thimble contributes,

$$\mathcal{Z} = \sum_a n_a e^{-i\text{Im}S[\phi_c^a]} \mathcal{Z}_a, \quad \mathcal{Z}_a = \int_{\mathcal{T}_a} \mathcal{D}\phi' e^{-\text{Re}S[\phi']},$$

and different constant phases can lead to cancellations

1. thimbles are clean analytically, but not necessarily optimal numerically
2. thimble decomposition can solve the sign problem only if a single thimble contributes (or all contributing thimbles have the same phase), *and* the Jacobian can be ignored

Method II: holomorphic flow

holomorphic (upward) flow equations

$$\frac{d\phi'}{dt} = \left(\frac{\partial S}{\partial \phi'} \right)^*$$

with initial conditions $\phi'(t=0) = \phi$

Method II: holomorphic flow

holomorphic (upward) flow equations

$$\frac{d\phi'}{dt} = \left(\frac{\partial S}{\partial \phi'} \right)^*$$

with initial conditions $\phi'(t=0) = \phi$

along the flow

$$\frac{d}{dt} \operatorname{Re} S = \frac{1}{2} \left[\frac{dS}{dt} + \left(\frac{dS}{dt} \right)^* \right] = \left| \frac{\partial S}{\partial \phi'} \right|^2 \geq 0$$

$$\frac{d}{dt} \operatorname{Im} S = \frac{1}{2i} \left[\frac{dS}{dt} - \left(\frac{dS}{dt} \right)^* \right] = \frac{1}{2i} \left[\left| \frac{\partial S}{\partial \phi'} \right|^2 - \left| \frac{\partial S}{\partial \phi'} \right|^2 \right] = 0$$

Method II: holomorphic flow

holomorphic (upward) flow equations

$$\frac{d\phi'}{dt} = \left(\frac{\partial S}{\partial \phi'} \right)^*$$

with initial conditions $\phi'(t=0) = \phi$

along the flow

$$\frac{d}{dt} \operatorname{Re} S = \frac{1}{2} \left[\frac{dS}{dt} + \left(\frac{dS}{dt} \right)^* \right] = \left| \frac{\partial S}{\partial \phi'} \right|^2 \geq 0$$

$$\frac{d}{dt} \operatorname{Im} S = \frac{1}{2i} \left[\frac{dS}{dt} - \left(\frac{dS}{dt} \right)^* \right] = \frac{1}{2i} \left[\left| \frac{\partial S}{\partial \phi'} \right|^2 - \left| \frac{\partial S}{\partial \phi'} \right|^2 \right] = 0$$

hence oscillations from e^{-S} are suppressed as t increases

Method II: holomorphic flow

holomorphic (upward) flow equations

$$\frac{d\phi'}{dt} = \left(\frac{\partial S}{\partial \phi'} \right)^*$$

with initial conditions $\phi'(t=0) = \phi$

along the flow

$$\frac{d}{dt} \operatorname{Re} S = \frac{1}{2} \left[\frac{dS}{dt} + \left(\frac{dS}{dt} \right)^* \right] = \left| \frac{\partial S}{\partial \phi'} \right|^2 \geq 0$$

$$\frac{d}{dt} \operatorname{Im} S = \frac{1}{2i} \left[\frac{dS}{dt} - \left(\frac{dS}{dt} \right)^* \right] = \frac{1}{2i} \left[\left| \frac{\partial S}{\partial \phi'} \right|^2 - \left| \frac{\partial S}{\partial \phi'} \right|^2 \right] = 0$$

hence oscillations from e^{-S} are suppressed as t increases

separate equation for the Jacobian

$$\frac{d\mathcal{J}}{dt} = (\mathcal{H}\mathcal{J})^*, \quad \mathcal{J}(t=0) = \mathbf{1}$$

Method II: holomorphic flow

two competing effects with increasing flow time t

Method II: holomorphic flow

two competing effects with increasing flow time t

- **Good:** $\text{Re}S$ grows \implies less oscillatory integrand

Method II: holomorphic flow

two competing effects with increasing flow time t

- **Good:** $\text{Re}S$ grows \implies less oscillatory integrand
- **Bad:** the Jacobian $\det \mathcal{J}(t, \phi)$ can become complex and can reintroduce cancellations

Method II: holomorphic flow

two competing effects with increasing flow time t

- **Good:** $\text{Re}S$ grows \implies less oscillatory integrand
- **Bad:** the Jacobian $\det \mathcal{J}(t, \phi)$ can become complex and can reintroduce cancellations

evolution of

$$\sigma(t) = \langle e^{i\theta} \rangle_{\text{PQ}}^{\text{def}}(t) = \frac{\mathcal{Z}}{\mathcal{Z}_{\text{PQ}}^{\text{def}}(t)}$$

is not necessarily monotonic!

Method III: numerical sign optimization

sign problem is weaker if

$$\sigma = \frac{\mathcal{Z}}{\mathcal{Z}_{\text{PQ}}^{\text{def}}} \quad \text{is maximal!}$$

Method III: numerical sign optimization

sign problem is weaker if

$$\sigma = \frac{\mathcal{Z}}{\mathcal{Z}_{\text{PQ}}^{\text{def}}} \quad \text{is maximal!}$$

parameterize the integration manifold via an ansatz depending a finite parameter set: $\mathcal{M} \rightarrow \mathcal{M}'(\{p\})$

Method III: numerical sign optimization

sign problem is weaker if

$$\sigma = \frac{\mathcal{Z}}{\mathcal{Z}_{\text{PQ}}^{\text{def}}} \quad \text{is maximal!}$$

parameterize the integration manifold via an ansatz depending a finite parameter set: $\mathcal{M} \rightarrow \mathcal{M}'(\{p\})$

optimize parameters $\{p\}$ by minimizing the cost function

$$\mathcal{F}(\{p\}) = -\log \sigma(\{p\}) = -\log \mathcal{Z} + \log \mathcal{Z}_{\text{PQ}}^{\text{def}}(\{p\})$$

Method III: numerical sign optimization

sign problem is weaker if

$$\sigma = \frac{\mathcal{Z}}{\mathcal{Z}_{\text{PQ}}^{\text{def}}} \quad \text{is maximal!}$$

parameterize the integration manifold via an ansatz depending a finite parameter set: $\mathcal{M} \rightarrow \mathcal{M}'(\{p\})$

optimize parameters $\{p\}$ by minimizing the cost function

$$\mathcal{F}(\{p\}) = -\log \sigma(\{p\}) = -\log \mathcal{Z} + \log \mathcal{Z}_{\text{PQ}}^{\text{def}}(\{p\})$$

i.e., by maximizing $\sigma(\{p\})$

Method III: numerical sign optimization

sign problem is weaker if

$$\sigma = \frac{\mathcal{Z}}{\mathcal{Z}_{\text{PQ}}^{\text{def}}} \quad \text{is maximal!}$$

parameterize the integration manifold via an ansatz depending a finite parameter set: $\mathcal{M} \rightarrow \mathcal{M}'(\{p\})$

optimize parameters $\{p\}$ by minimizing the cost function

$$\mathcal{F}(\{p\}) = -\log \sigma(\{p\}) = -\log \mathcal{Z} + \log \mathcal{Z}_{\text{PQ}}^{\text{def}}(\{p\})$$

i.e., by maximizing $\sigma(\{p\})$

since \mathcal{Z} is contour-invariant (within homology), in practice:

$$\min_{\{p\}} \mathcal{F}(\{p\}) \iff \min_{\{p\}} \log \mathcal{Z}_{\text{PQ}}^{\text{def}}(\{p\})$$

Method III: numerical sign optimization

sign problem is weaker if

$$\sigma = \frac{\mathcal{Z}}{\mathcal{Z}_{\text{PQ}}^{\text{def}}} \quad \text{is maximal!}$$

parameterize the integration manifold via an ansatz depending a finite parameter set: $\mathcal{M} \rightarrow \mathcal{M}'(\{p\})$

optimize parameters $\{p\}$ by minimizing the cost function

$$\mathcal{F}(\{p\}) = -\log \sigma(\{p\}) = -\log \mathcal{Z} + \log \mathcal{Z}_{\text{PQ}}^{\text{def}}(\{p\})$$

i.e., by maximizing $\sigma(\{p\})$

since \mathcal{Z} is contour-invariant (within homology), in practice:

$$\min_{\{p\}} \mathcal{F}(\{p\}) \iff \min_{\{p\}} \log \mathcal{Z}_{\text{PQ}}^{\text{def}}(\{p\})$$

gradient of the cost function

$$\nabla_p \mathcal{F}(\{p\}) = \nabla_p \log \mathcal{Z}_{\text{PQ}}^{\text{def}}(\{p\}) = -\langle \nabla_p \text{ReS}_{\text{PQ}}^{\text{eff}}(\{p\}) \rangle_{\text{PQ}}^{\text{def}}$$

Method III: numerical sign optimization – ansatz

in 1D, parameterize a smooth, single-valued deformation

$$z(\phi) = \phi + if(\phi, \{p\}), \quad f : \mathbb{R} \rightarrow \mathbb{R}$$

Method III: numerical sign optimization – ansatz

in 1D, parameterize a smooth, single-valued deformation

$$z(\phi) = \phi + i f(\phi, \{p\}), \quad f : \mathbb{R} \rightarrow \mathbb{R}$$

choose nodes x_i , $i = 1, 2, \dots, M$, with f linearly interpolating between them,

$$f(\phi, x_{j(\phi)}, x_{j(\phi)+1}, y_{j(\phi)}, y_{j(\phi)+1}) = \frac{y_{j(\phi)}(x_{j(\phi)+1} - \phi)}{x_{j(\phi)+1} - x_{j(\phi)}} + \frac{y_{j(\phi)+1}(\phi - x_{j(\phi)})}{x_{j(\phi)+1} - x_{j(\phi)}}$$

Method III: numerical sign optimization – ansatz

in 1D, parameterize a smooth, single-valued deformation

$$z(\phi) = \phi + i f(\phi, \{p\}), \quad f : \mathbb{R} \rightarrow \mathbb{R}$$

choose nodes x_i , $i = 1, 2, \dots, M$, with f linearly interpolating between them,

$$f(\phi, x_{j(\phi)}, x_{j(\phi)+1}, y_{j(\phi)}, y_{j(\phi)+1}) = \frac{y_{j(\phi)}(x_{j(\phi)+1} - \phi)}{x_{j(\phi)+1} - x_{j(\phi)}} + \frac{y_{j(\phi)+1}(\phi - x_{j(\phi)})}{x_{j(\phi)+1} - x_{j(\phi)}}$$

- parameters are simply $\{p\} = \{y_0, \dots, y_M\}$

Method III: numerical sign optimization – ansatz

in 1D, parameterize a smooth, single-valued deformation

$$z(\phi) = \phi + if(\phi, \{p\}), \quad f : \mathbb{R} \rightarrow \mathbb{R}$$

choose nodes x_i , $i = 1, 2, \dots, M$, with f linearly interpolating between them,

$$f(\phi, x_{j(\phi)}, x_{j(\phi)+1}, y_{j(\phi)}, y_{j(\phi)+1}) = \frac{y_{j(\phi)}(x_{j(\phi)+1} - \phi)}{x_{j(\phi)+1} - x_{j(\phi)}} + \frac{y_{j(\phi)+1}(\phi - x_{j(\phi)})}{x_{j(\phi)+1} - x_{j(\phi)}}$$

- parameters are simply $\{p\} = \{y_0, \dots, y_M\}$
- Jacobian is cheap

$$\mathcal{J}(\phi) = 1 + i \frac{y_{j(\phi)+1} - y_{j(\phi)}}{\Delta}$$

Method III: numerical sign optimization – ansatz

in 1D, parameterize a smooth, single-valued deformation

$$z(\phi) = \phi + if(\phi, \{p\}), \quad f : \mathbb{R} \rightarrow \mathbb{R}$$

choose nodes x_i , $i = 1, 2, \dots, M$, with f linearly interpolating between them,

$$f(\phi, x_{j(\phi)}, x_{j(\phi)+1}, y_{j(\phi)}, y_{j(\phi)+1}) = \frac{y_{j(\phi)}(x_{j(\phi)+1} - \phi)}{x_{j(\phi)+1} - x_{j(\phi)}} + \frac{y_{j(\phi)+1}(\phi - x_{j(\phi)})}{x_{j(\phi)+1} - x_{j(\phi)}}$$

- parameters are simply $\{p\} = \{y_0, \dots, y_M\}$
- Jacobian is cheap

$$\mathcal{J}(\phi) = 1 + i \frac{y_{j(\phi)+1} - y_{j(\phi)}}{\Delta}$$

- index function $j(\phi) = \text{floor}[(\phi - x_1)/\Delta]$ with $\Delta = x_{i+1} - x_i (= \text{const.})$

Method III: numerical sign optimization – ansatz

in 1D, parameterize a smooth, single-valued deformation

$$z(\phi) = \phi + if(\phi, \{p\}), \quad f : \mathbb{R} \rightarrow \mathbb{R}$$

choose nodes x_i , $i = 1, 2, \dots, M$, with f linearly interpolating between them,

$$f(\phi, x_{j(\phi)}, x_{j(\phi)+1}, y_{j(\phi)}, y_{j(\phi)+1}) = \frac{y_{j(\phi)}(x_{j(\phi)+1} - \phi)}{x_{j(\phi)+1} - x_{j(\phi)}} + \frac{y_{j(\phi)+1}(\phi - x_{j(\phi)})}{x_{j(\phi)+1} - x_{j(\phi)}}$$

- parameters are simply $\{p\} = \{y_0, \dots, y_M\}$
- Jacobian is cheap

$$\mathcal{J}(\phi) = 1 + i \frac{y_{j(\phi)+1} - y_{j(\phi)}}{\Delta}$$

- index function $j(\phi) = \text{floor}[(\phi - x_1)/\Delta]$ with $\Delta = x_{i+1} - x_i (= \text{const.})$
- optimization can be done with standard gradient-based methods

Method III: numerical sign optimization – ansatz

in 1D, parameterize a smooth, single-valued deformation

$$z(\phi) = \phi + i f(\phi, \{p\}), \quad f : \mathbb{R} \rightarrow \mathbb{R}$$

choose nodes $x_i, i = 1, 2, \dots, M$, with f linearly interpolating between them,

$$f(\phi, x_{j(\phi)}, x_{j(\phi)+1}, y_{j(\phi)}, y_{j(\phi)+1}) = \frac{y_{j(\phi)}(x_{j(\phi)+1} - \phi)}{x_{j(\phi)+1} - x_{j(\phi)}} + \frac{y_{j(\phi)+1}(\phi - x_{j(\phi)})}{x_{j(\phi)+1} - x_{j(\phi)}}$$

- parameters are simply $\{p\} = \{y_0, \dots, y_M\}$
- Jacobian is cheap

$$\mathcal{J}(\phi) = 1 + i \frac{y_{j(\phi)+1} - y_{j(\phi)}}{\Delta}$$

- index function $j(\phi) = \text{floor}[(\phi - x_1)/\Delta]$ with $\Delta = x_{i+1} - x_i (= \text{const.})$
- optimization can be done with standard gradient-based methods
- by increasing M and decreasing Δ one can approximate any continuous function, so piecewise linearity is not an assumption, however continuity and being single-valued are

Toy models - 1D integrals

we benchmark all three methods on four fermionic toy partition functions, covering

Toy models - 1D integrals

we benchmark all three methods on four fermionic toy partition functions, covering

- **noncompact** auxiliary-field integrals (Hubbard-/Gross-Neveu-like)

Toy models - 1D integrals

we benchmark all three methods on four fermionic toy partition functions, covering

- **noncompact** auxiliary-field integrals (Hubbard-/Gross-Neveu-like)
- **compact** variable integrals (Thirring-/Chern-Simons-like)

Toy models - 1D integrals

we benchmark all three methods on four fermionic toy partition functions, covering

- **noncompact** auxiliary-field integrals (Hubbard-/Gross-Neveu-like)
- **compact** variable integrals (Thirring-/Chern-Simons-like)

we monitor improvement with the average phase

$$\sigma = \frac{\mathcal{Z}}{\mathcal{Z}_{\text{PQ}}^{\text{def}}}$$

Model 1: One-site Hubbard-like integral

$$\mathcal{Z}_H = \sqrt{\frac{1}{2\pi U}} \int_{-\infty}^{\infty} d\phi \left(1 + e^{\mu+U/2+i\phi}\right)^2 e^{-\phi^2/(2U)}$$

- integration domain: $\phi \in \mathbb{R}$ (noncompact)
- infinitely many critical points ϕ_c
- (dimensionless) parameters: chemical potential μ and interaction U
- sign problem away from half filling ($\mu/U = 1/2$)

Model 2: Gross-Neveu-like (0+1D) integral

$$\mathcal{Z}_{\text{GN}} = \sqrt{\frac{N}{\pi G}} \int_{-\infty}^{\infty} d\phi \left(p^2 + (m + \phi)^2 \right)^N \exp\left(-\frac{N}{G}\phi^2\right)$$

- integration domain $\phi \in \mathbb{R}$ (noncompact)
- finite number of critical points
- parameters: number of flavors N , coupling G , momentum p and mass m
- sign problem induced by complex coupling G

Model 3: heavy-dense Thirring-like integral

$$\mathcal{Z}_T = \int_{-\pi}^{\pi} d\phi \exp\left(\frac{N}{g^2} \cos \phi + i\phi + \mu\right)$$

- integration domain $\phi \in [-\pi, \pi]$ (compact),
- single critical point at $z_c = i \operatorname{arsinh}(g^2/N)$
- action becomes singular at $\operatorname{Re}z = \pm\pi$
- parameters: number of flavors N , coupling g and chemical potential μ

Model 4: Chern–Simons–like with fermions (compact)

$$\mathcal{Z}_{\text{CS}} = \frac{e^{-ik\pi}}{2^N} \int_{-\pi}^{\pi} d\phi (\cosh m - \cos \phi)^N e^{ik\phi}$$

- integration domain $\phi \in [-\pi, \pi]$ (compact)
- interesting case: $k < N$ (for $k > N$ $\mathcal{Z} = 0$ and $k = N$ critical points are imaginary infinity)
- parameters: number of flavors N , topological charge k and mass m
- sign problem from topological term $e^{ik\phi}$

Summary of toy models

model	Hubbard	Gross-Neveu	Thirring	Chern-Simons
\int domain	noncompact	noncompact	compact	compact
# of CPs	infinite	finite	1	finite
sign problem	μ	complex G	μ	θ term
Im-thimble	finite	finite	divergent	finite

Figure: Flow time evolution (Hubbard)

$$\mathcal{Z}_H = \sqrt{\frac{1}{2\pi U}} \int_{-\infty}^{\infty} d\phi \left(1 + e^{\mu+U/2+i\phi}\right)^2 e^{-\phi^2/(2U)}$$

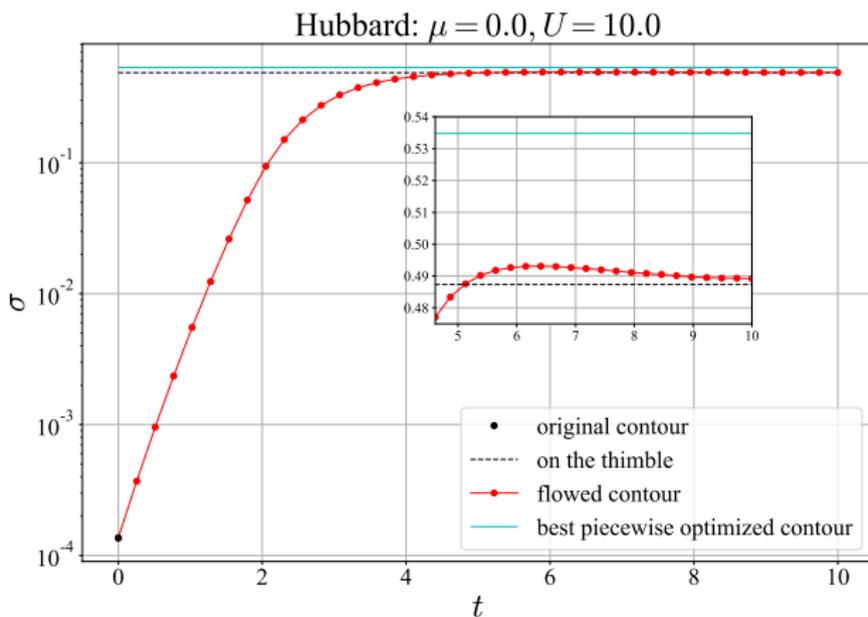


Figure: Contours in the complex plane (Hubbard)

$$\mathcal{Z}_H = \sqrt{\frac{1}{2\pi U}} \int_{-\infty}^{\infty} d\phi \left(1 + e^{\mu+U/2+i\phi}\right)^2 e^{-\phi^2/(2U)}$$

Hubbard: $\mu = 0.0, U = 10.0$

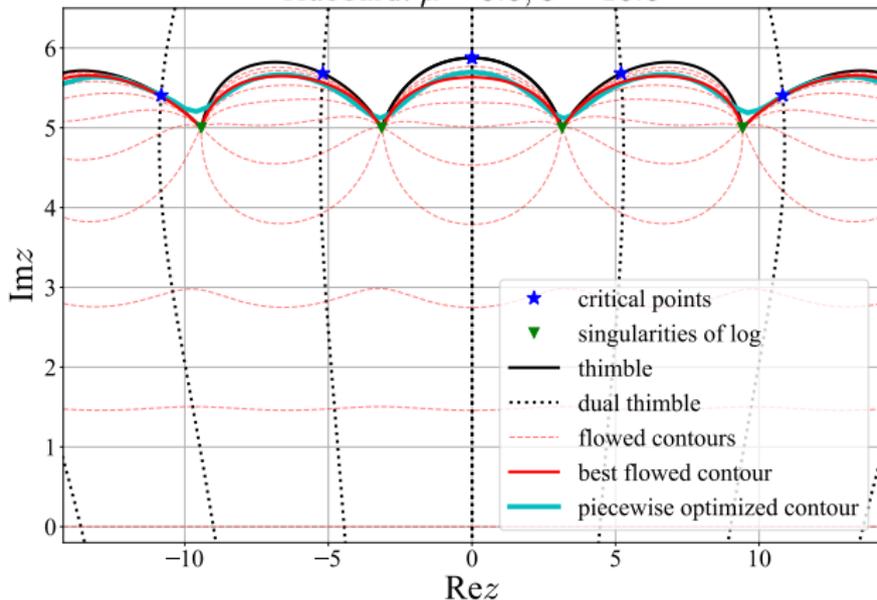


Figure: Flow time evolution (Gross-Neveu)

$$\mathcal{Z}_{\text{GN}} = \sqrt{\frac{N}{\pi G}} \int_{-\infty}^{\infty} d\phi \left(p^2 + (m + \phi)^2 \right)^N \exp\left(-\frac{N}{G} \phi^2 \right),$$

Gross-Neveu: $N = 2, p = 1.0, m = 0.0, G = 1.75 + i$

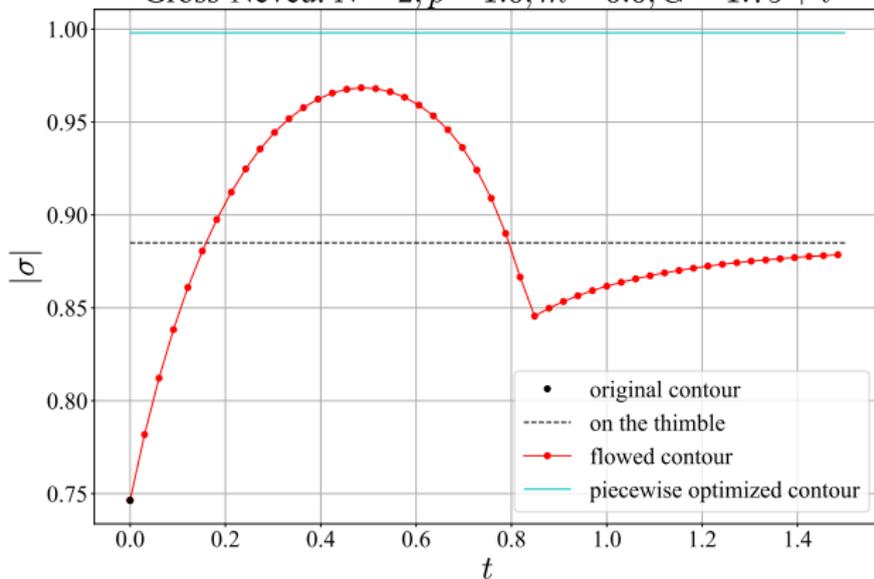


Figure: Contours in the complex plane (Gross-Neveu)

$$\mathcal{Z}_{\text{GN}} = \sqrt{\frac{N}{\pi G}} \int_{-\infty}^{\infty} d\phi \left(p^2 + (m + \phi)^2 \right)^N \exp\left(-\frac{N}{G} \phi^2\right),$$

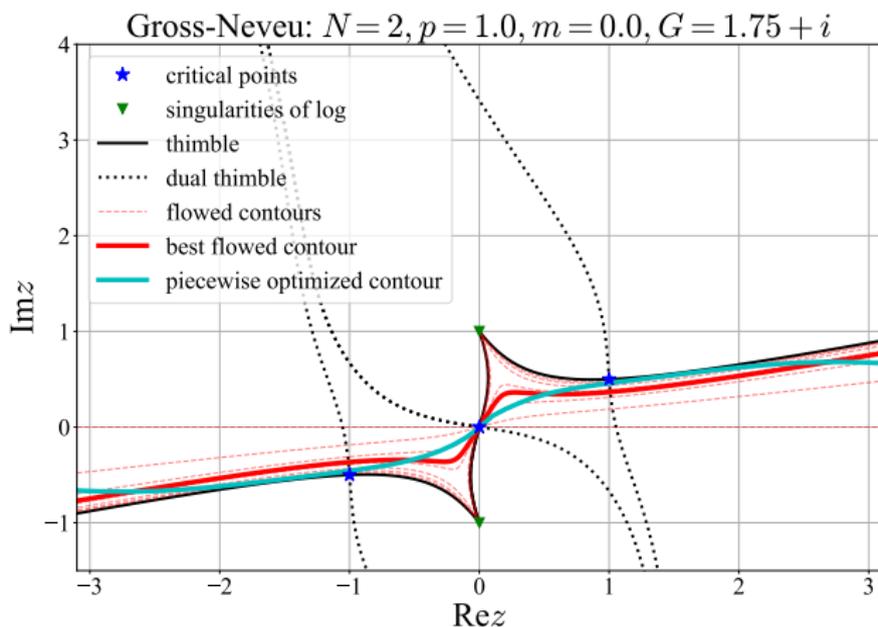


Figure: Flow time evolution (Thirring)

$$\mathcal{Z}_T = \int_{-\pi}^{\pi} d\phi \exp\left(\frac{N}{g^2} \cos \phi + i\phi + \mu\right)$$

Thirring: $N = 2, g = 1.08$

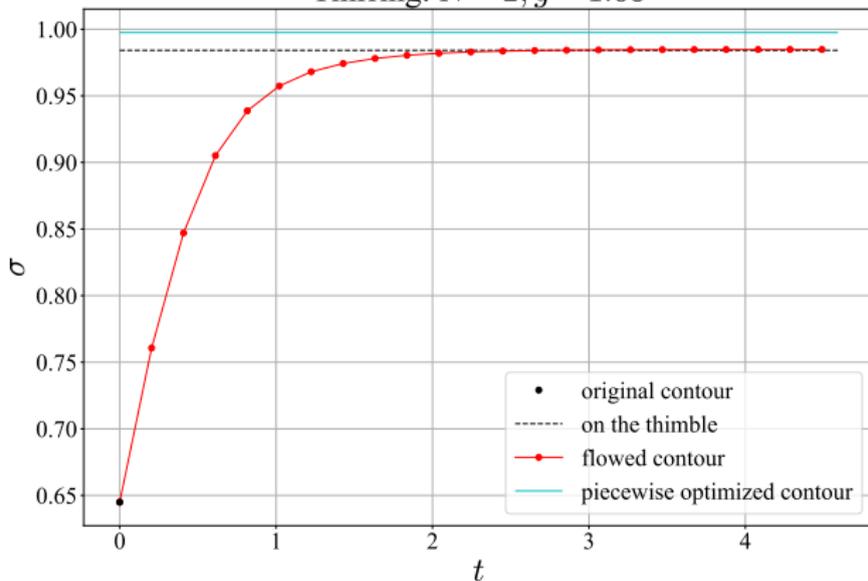


Figure: Contours in the complex plane (Thirring)

$$\mathcal{Z}_T = \int_{-\pi}^{\pi} d\phi \exp\left(\frac{N}{g^2} \cos \phi + i\phi + \mu\right)$$

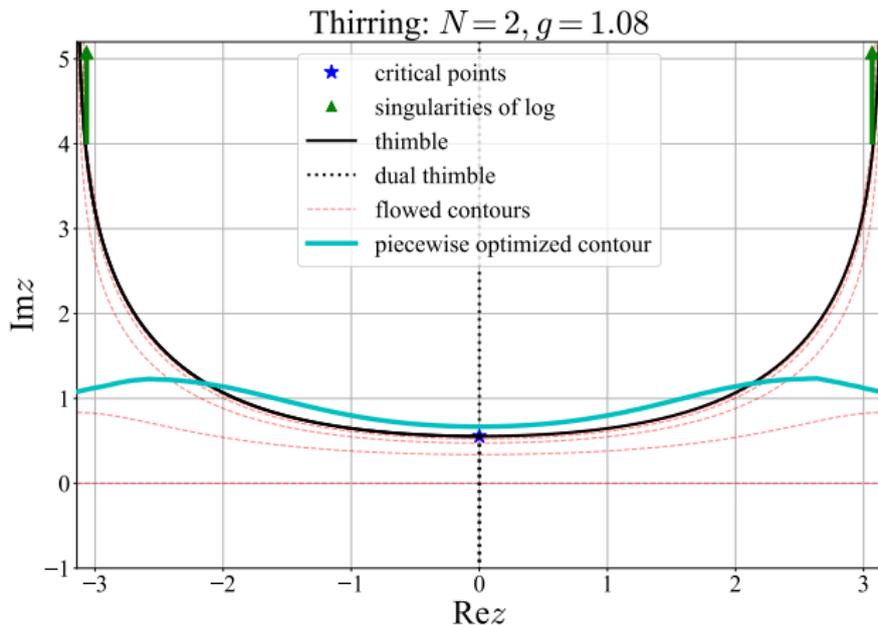


Figure: Flow time evolution (Chern-Simons)

$$\mathcal{Z}_{\text{CS}} = \frac{e^{-ik\pi}}{2^N} \int_{-\pi}^{\pi} d\phi (\cosh m - \cos \phi)^N e^{ik\phi}$$

Chern-Simons: $N = 2, m = 0.4, k = 1$

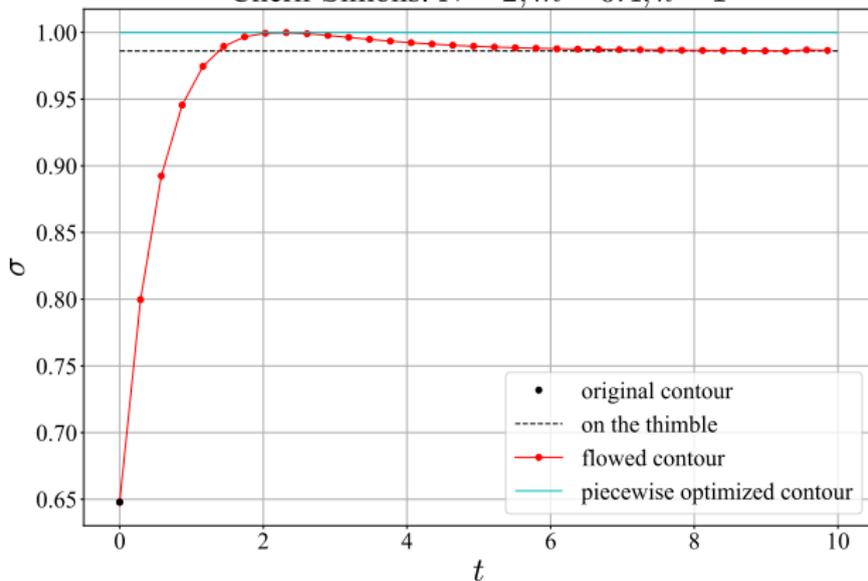
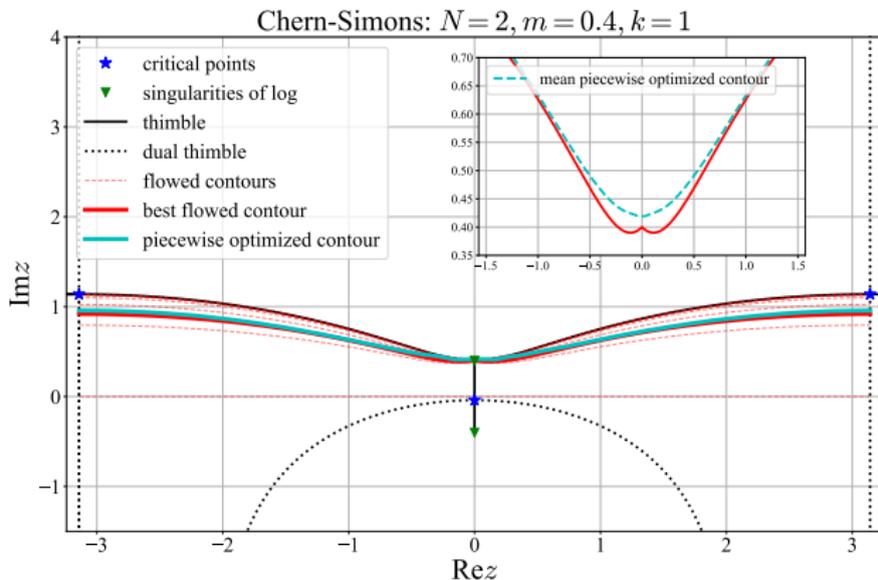


Figure: Contours in the complex plane (Chern-Simons)

$$\mathcal{Z}_{\text{CS}} = \frac{e^{-ik\pi}}{2^N} \int_{-\pi}^{\pi} d\phi (\cosh m - \cos \phi)^N e^{ik\phi}$$



Summary and outlook

across the toy models

Summary and outlook

across the toy models

- Lefschetz thimbles are not the contours that maximize σ

Summary and outlook

across the toy models

- Lefschetz thimbles are not the contours that maximize σ
- holomorphic flow often has an optimal finite t^*

Summary and outlook

across the toy models

- Lefschetz thimbles are not the contours that maximize σ
- holomorphic flow often has an optimal finite t^*
- smooth single-valued optimized contours can yield the best σ among tested methods

Summary and outlook

across the toy models

- Lefschetz thimbles are not the contours that maximize σ
- holomorphic flow often has an optimal finite t^*
- smooth single-valued optimized contours can yield the best σ among tested methods
- natural limitation: one-dimensional integrals only

Summary and outlook

across the toy models

- Lefschetz thimbles are not the contours that maximize σ
- holomorphic flow often has an optimal finite t^*
- smooth single-valued optimized contours can yield the best σ among tested methods
- natural limitation: one-dimensional integrals only
- for not-too-large dimensional integrals this interpolant ansatz can be generalized via simplices

Summary and outlook

across the toy models

- Lefschetz thimbles are not the contours that maximize σ
- holomorphic flow often has an optimal finite t^*
- smooth single-valued optimized contours can yield the best σ among tested methods
- natural limitation: one-dimensional integrals only
- for not-too-large dimensional integrals this interpolant ansatz can be generalized via simplices
- higher dimensional integrals have further important features (spacetime and internal symmetries) that could be used in constructing a good ansatz for optimization

The End

Thank you!

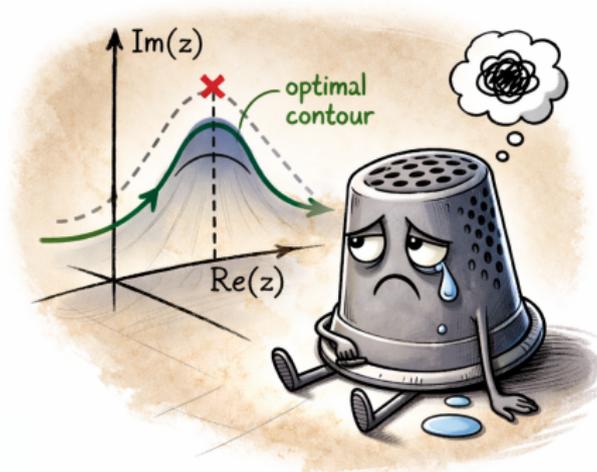


Figure: via ChatGPT