# Particle physics 

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## Cross section

Operative definition

$$
\sigma=\frac{N_{\text {events }}}{N_{t} \frac{N_{b}}{A_{b}}}
$$

Uniformly distributed beam, crossing the whole target in time $t$

$$
N_{b}=\frac{N_{b}}{t} t=\frac{\Delta N_{b}}{\Delta t} \Delta t=\rho_{b} A_{b} \frac{\Delta x}{\Delta t} \Delta t=\rho_{b} A_{b} v \Delta t=\Phi A_{b} \Delta t
$$

$\rho_{b}$ : beam density (n. particle/unit volume)
$\Phi=\rho_{b} v$ : beam flux ( n . particle crossing unit area $\perp$ beam per unit time)
Constant event rate $N_{\text {events }}=\frac{N_{\text {events }}}{t} t=\frac{\Delta N_{\text {events }}}{\Delta t} \Delta t$

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\sigma=\frac{\frac{\Delta N_{\text {events }}}{\Delta t}}{N_{t} \frac{1}{A_{b}} \frac{\Delta N_{b}}{\Delta t}}=\frac{\Delta N_{\text {events }}}{\Delta t N_{t} \Phi}
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Cross section $=\mathrm{n}$. scattering events per unit time, unit target, unit flux

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## Scattering processes

Elastic scattering

$$
a b \longrightarrow a b
$$

- same type (and number) of particles in and out
- momenta and spin component can change

Inelastic scattering

$$
a b \longrightarrow X_{1} X_{2} \ldots X_{n}
$$

- different particles in and out
- kinematical and dynamical constraints restrict the allowed inelastic processeses


## Resonances

Distinctive sign of unstable particle being created as intermediate state: peak in the cross section as a function of energy (resonance)

- Position of the peak $\rightarrow$ mass $m$
- Width of the peak $\rightarrow$ decay width $\Gamma=1$ /lifetime "Hand-waving" argument: wave-function of unstable system of $E \approx m$ decaying exponentially in time with lifetime $1 / \Gamma$

$$
\begin{aligned}
\psi(t) & =\psi(0) e^{-i m t} e^{-\frac{\Gamma}{2} t} \longrightarrow|\psi(t)|^{2}=|\psi(0)| e^{-\Gamma t} \\
\tilde{\psi}(E) & =\int d t e^{i E t} \psi(t)=\frac{i \psi(0)}{E-m+i \frac{\Gamma}{2}}
\end{aligned}
$$

If the unstable system is formed in a scattering experiment at energy $E$, $\sigma(E)$ near $m \propto$ probability of observing the unstable system with energy $E$

$$
\sigma(E) \propto|\tilde{\psi}(E)|^{2}=\frac{|\psi(0)|^{2}}{(E-m)^{+}\left(\frac{\Gamma}{2}\right)^{2}}=\sigma_{\max } \frac{\left(\frac{\Gamma}{2}\right)^{2}}{(E-m)^{2}+\left(\frac{\Gamma}{2}\right)^{2}}
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$\phi$ meson $(s \bar{s})$ in lead-lead collisions
$m_{\phi}=1019.461 \pm 0.016 \mathrm{MeV}, \Gamma_{\phi}=4.249 \pm 0.013 \mathrm{MeV}[4]$

## Describing interactions: Feynman diagrams

Interactions can be described as exchange of particles
E.g.: $e^{-} e^{+}$scattering mediated by $\mathrm{EM}=$ exchange of (one or more) $\gamma$

- particle 1 emits/absorbs mediator absorbed/emitted by particle 2...
- ... or the other way around?
- both, and none of the two: not well defined who does what on such short time-scales

What matters is the exchange, not who emits/absorbs

> (Also: do not take this picture too literally)

Fundamental processes: emission/absorption of an interaction particle from matter particle or from another interaction particle (vertex) Quantities conserved at vertex $\rightarrow$ automatically conserved by interaction: energy/momentum, angular momentum, electric charge...

## Electromagnetic interactions

Interaction vertex in Quantum ElectroDynamics (QED)

- electron enters, emits/absorbs photon, exits (time flows upwards)
- only vertex
- same for any other negatively charged lepton, or for quarks
- same for antiparticles, except arrow is drawn reversed (time still flows upwards)

Diagrams like these are known as Feynman diagrams: more than pictorial representation of a process (in due time)

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## Strong interactions

Interaction vertices in Quantum ChromoDynamics (QCD)

- quark enters, exchanges gluon, exits
- similar to QED, but quarks and gluons carry also colour
- colour of $q$ can change but overall conserved at vertex ("difference" carried by $g$ )
- 3 quark colours, 8 gluon types
( $3 \times 3=9$ combinations, but the one leaving all colours unchanged is absent)
- also vertices involving only 3 or 4 gluons (gluons self-interact)


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## Weak interactions

Interaction vertices for weak interactions

- charged current: negatively charged lepton enters, emits $W^{-}$/absorbs $W^{+}$and turns into neutrino...
- ... or neutrino enters, emits $W^{+}$/absorbs $W^{-}$and turns into neg. charged lepton
- similarly with antiparticles
- leptons from same family are involved: $\left(I^{-}, \nu_{l}\right),\left(I^{+}, \bar{\nu}_{l}\right)$
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## Weak interactions of quarks

N . current same as with leptons, ch. current analogue $\left(e^{-}, \nu_{e}\right) \rightarrow(u, d)$ Instead $\left(e^{-}, \nu_{e}\right) \rightarrow\left(u, d^{\prime}\right)$ with $d^{\prime}$ a superposition of $d$ and $s$ quarks

$$
\left|d^{\prime}\right\rangle=\cos \theta_{C}|d\rangle+\sin \theta_{C}|s\rangle
$$

Needed to explain $K \rightarrow$ hadrons, where $s / \bar{s}$ turns into $d / \bar{d}$
Better: $\left(\ell^{-}, \nu_{\ell}\right) \rightarrow\left(u, d^{\prime}\right),\left(c, s^{\prime}\right),\left(t, b^{\prime}\right)$ with $d^{\prime}, s^{\prime}, b^{\prime}$ lin. sup. of $d, s, b$ Unitary matrix of mixing coefficients is the Cabibbo-Kobayashi-Maskawa (CKM) matrix

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## Diagrams with antiparticles

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Vertices can be "rotated" to put fermion and antifermion on the same side of the process

## Conservation laws

From interaction vertices one can read off conservation laws

|  | EM | strong | weak |
| :---: | :---: | :---: | :---: |
| electric charge | yes | yes | yes |
| lepton type | yes | - | no |
| flavour (=quark type) | yes | yes | no |
| lepton family | yes | - | yes (if massless) |
| quark family | yes | yes | no |
| lepton number | yes | - | yes |
| quark number | yes | yes | yes |

Lepton type/flavour: $n_{f}-n_{\bar{f}}$
Lepton family number: $L_{\ell}=\left(n_{\ell^{-}}-n_{\ell^{+}}\right)+\left(n_{\nu_{\ell}}-n_{\bar{\nu}_{\ell}}\right)$
"Quark family": $\left(n_{u}-n_{\bar{u}}\right)+\left(n_{d}-n_{\bar{d}}\right)$, etc.
(never used: either more detailed cons. law exists, or not conserved due to quark mixing)
Lepton number: $L=\sum_{\ell=e, \mu, \tau} L_{\ell}$
Quark number: $\mathcal{Q}=\sum_{q} n_{q}-n_{\bar{q}}$

## Conservation laws (contd.)

Flavour numbers:

$$
\begin{array}{lll}
U=n_{u}-n_{\bar{u}} & C=n_{c}-n_{\bar{c}} & T=n_{t}-n_{\bar{t}} \\
D=-n_{d}+n_{\bar{d}} & S=-n_{s}+n_{\bar{s}} & B=-n_{b}+n_{\bar{b}}
\end{array}
$$

Quark number is equivalent to baryon number

$$
\begin{aligned}
& \sum_{q} n_{q}=3 n_{\text {baryons }}+n_{\text {mesons }} \\
& \sum_{q} n_{\bar{q}}=3 n_{\text {antibaryons }}+n_{\text {mesons }} \\
& \Longrightarrow \mathcal{Q}=3\left(n_{\text {baryons }}-n_{\text {antibaryons }}\right)=3 \mathcal{B}
\end{aligned}
$$

If flavour is conserved, $U$ and $D$ are traded for $\mathcal{B}$ and electric charge $Q$

$$
\begin{aligned}
& \mathcal{B}=\frac{1}{3}(U+C+T-D-S-B) \\
& Q=\frac{2}{3}(U+C+T)+\frac{1}{3}(D+S+B)
\end{aligned}
$$

If an interaction conserves a certain particle number it cannot be responsible for decays in which this number is violated (e.g., strangeness changing processes cannot be due to strong int.)

## Basic processes

- Vertex diagrams describe how the interaction works at the most fundamental level but cannot represent a true physical process due to lack of energy-momentum conservation,
- Actual physical process described combining two or more vertices

- energy and momentum are conserved, but exchanged photon not on mass shell, $p_{\gamma}^{2} \neq 0=m_{\gamma}^{2}$
- internal lines represent virtual particles, not real particles, so not on shell
- external lines represent real particles, must be on-shell $p_{i}^{2}=m_{i}^{2}$


## Basic processes in QED



Møller scattering

$$
e^{-} e^{-} \rightarrow e^{-} e^{-}
$$

electrons are indistinguishable, cannot say which one is going left/right after photon exchange $\rightarrow$ must take both possibilities into account

## Basic processes in QED



Bhabha scattering

$$
e^{-} e^{+} \rightarrow e^{-} e^{+}
$$

$e^{-} e^{+}$can exchange a photon or annihilate into a photon $\rightarrow$ must take both possibilities into account

## Basic processes in QED



Three more basic QED processes:

- electron-positron annihilation $e^{-} e^{+} \rightarrow \gamma \gamma$
- electron-positron pair creation $\gamma \gamma \rightarrow e^{-} e^{+}$
- Compton scattering $\gamma e^{-} \rightarrow \gamma e^{-}$
(For each of these processes there is a second diagram: can you draw it?)


## What matters in a Feynman diagram

Feynman diagrams are not an accurate depiction of particle trajectories!


- "time" not well defined, except "before" and "after" of the process
- "when" annihiliation/creation happens are meaningless questions
- all that matters is the topology of the diagram, i.e., connectivity properties (once initial and final states are given)


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## Higher order diagrams

Infinity of ways in which one can combine vertices

Left: light-by-light (Delbrück) scattering, Right: $e^{-} e^{-}$elastic scattering

- Diagrams have different "weights", i.e., are more or less important
- Each vertex contributes a factor $e$ to the weight of a diagram
- two vertices $\rightarrow \propto \alpha, \alpha=e^{2} /(4 \pi) \simeq 1 / 137$ fine structure constant
- four vertices $\rightarrow \alpha^{2}$, relatively suppressed
- Describing process to given precision requires limited n. of diagrams
- Vertex weighting factor: coupling constant

1 EM (e), 1 strong, 2 weak both $\propto e$ via Weinberg angle (EW unification)

## Basic charged weak current processes



Muon decay $\mu^{-} \rightarrow e^{-} \bar{\nu}_{e} \nu_{\mu}$
Neutron $\beta$-decay $n \rightarrow p e^{-} \bar{\nu}_{e}$ (same diagram plus spectator quarks)
Pion decay: $\pi^{-} \rightarrow \mu^{-} \bar{\nu}_{\mu}$ (replace e with $\mu$ )
Strangeness-changing processes: $\Lambda^{0}(u d s) \rightarrow p(u u d) \pi^{-}(d \bar{u}) \quad(\Delta S=1)$

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## References

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