

# HOMWORK N. 4

## Gauge theories

1. The *Wilson line*  $W_{\mathcal{C}}$  associated with a path  $\mathcal{C}$  connecting  $x_0$  and  $x_1$ ,

$$\mathcal{C} : [0, T] \rightarrow \mathbb{R}^4, \quad s \rightarrow x^\mu(s),$$

$x(0) = x_0$  and  $x(T) = x_1$ , is defined as follows

$$W_{\mathcal{C}} \equiv \text{Texp} \left\{ -ig \int_{\mathcal{C}} A_\mu(X) dX^\mu \right\} = \text{Texp} \left\{ -ig \int_0^T ds A_\mu(x(s)) \dot{x}^\mu(s) \right\},$$

where  $A_\mu = A_\mu^a t^a$  with  $t^a = t^{a\dagger}$  Hermitean matrices providing a representation of the Lie algebra associated to the gauge group and  $A_\mu^a$  the gauge fields;  $\dot{x}^\mu = dx^\mu/ds$  the four-vector tangent to the path, and  $\text{Texp}$  the time-ordered exponential,

$$\begin{aligned} \text{Texp} \left\{ \int_{s_0}^{s_1} ds f(s) \right\} &= \sum_{n=0}^{\infty} \int_{s_0}^{s_1} ds_1 \int_{s_0}^{s_1} ds_2 \dots \int_{s_0}^{s_1} ds_{n-1} \int_{s_0}^{s_1} ds_n \\ &\quad \times \theta(s_1 - s_2) \dots \theta(s_{n-1} - s_n) f(s_1) f(s_2) \dots f(s_{n-1}) f(s_n). \end{aligned}$$

where it is understood that the  $n = 0$  term is just  $\mathbf{1}$  (the identity matrix in the appropriate number of dimensions).

1.1. Let  $W_{\mathcal{C}}(t)$  denote the Wilson line along  $\mathcal{C}$  but only up to  $x(t)$ ,  $t < T$ ,

$$W_{\mathcal{C}}(t) \equiv \text{Texp} \left\{ -ig \int_0^t ds A_\mu(x(s)) \dot{x}^\mu(s) \right\}.$$

Show that  $W_{\mathcal{C}}(t)$  obeys the following differential equation,

$$\frac{dW_{\mathcal{C}}(t)}{dt} = -ig A_\mu(x(t)) \dot{x}^\mu(t) W_{\mathcal{C}}(t), \quad (*)$$

with initial condition  $W_{\mathcal{C}}(0) = \mathbf{1}$ . (Hint: use the theta functions in the definition of the time-ordered exponential to explicitly limit the integration ranges over  $s_1, \dots, s_n$ .)

1.2. From Equation (\*) obtain a differential equation for  $W_{\mathcal{C}}(t)^\dagger$ . From the relation  $W_{\mathcal{C}}(t)W_{\mathcal{C}}(t)^{-1} = \mathbf{1}$  derive a differential equation for  $W_{\mathcal{C}}(t)^{-1}$ . Compare the two equations and the corresponding initial conditions, and use uniqueness of the solution of the Cauchy problem to conclude that  $W_{\mathcal{C}}(t)$  is a unitary matrix.

1.3. Under a finite gauge transformation  $U(x) = e^{i\alpha_a(x)t^a}$ , the gauge field  $A_\mu(x)$  transforms as

$$A_\mu^U(x) = U(x)A_\mu(x)U(x)^\dagger + \frac{1}{ig}U(x)\partial_\mu U(x)^\dagger.$$

Imposing the axial gauge condition on  $A_\mu^U$ , i.e.,  $A_3^U = 0$ , derive a differential equation for  $U(x)$ , and write its solution in terms of a suitable (family of) Wilson lines.

2. Consider Faddeev-Popov quantisation of a gauge theory using the gauge-fixing functional

$$f^a[A; x] = n^\mu A_\mu^a.$$

2.1. Derive the Faddeev-Popov determinant corresponding to the above choice of  $f^a[A; x]$ , and its representation as a Grassmann integral over ghost and antighost fields.

2.2. Show that for the choice  $B[f] = \prod_x \delta(f[A; x])$  for the functional implementing the gauge condition in the path integral, the ghosts decouple from the gauge fields.

2.3. Derive the gauge-field propagator in momentum space for the choice  $B[f] = e^{\frac{i}{2\xi} \int d^4x f^a f^a}$ . (Hint: the kernel of the quadratic part of the Lagrangean leaves invariant both the subspace spanned by the momentum  $p^\mu$  and by  $n^\mu$  and its orthogonal complement, and so must do its inverse. What does this tell you about the most general form of the propagator?)