

STATUS OF THE SUPERWEAK EXTENSION OF THE STANDARD MODEL AND MUON $g - 2$

based on

arXiv:1812.11189 (*Symmetry*), 1911.07082 (*PRD*), 2104.11248 (*JCAP*), 2104.14571 (*PRD*), 2105.13360 (*J.Phys.G*), 2204.07100 (*PRD*), 2301.07961 (*JHEP*), 2301.06621 (*PRD*), 2305.11931 (*PRDL*) and correspondence with BMW collaboration
with S. Iwamoto, T.J. Kärkkäinen, I. Nándori, Z. Péli, K. Seller, Zs. Szép

OUTLINE

1. **Motivation**: status of particle physics
 - Colliders
 - Cosmology
 - **Muon anomalous magnetic moment**
2. Superweak $U(1)_Z$ extension of SM (**SWSM**)
3. Neutrino masses and dark matter candidate
4. Vacuum stability and scalar sector constraints
5. Contribution to M_W
6. Conclusions
7. *Appendix*: Constraints from non-standard interactions

MOTTO

*Rough estimates of BSM effects
can easily be deceptive*

example: discovery of the Higgs particle came much faster than expected at the time of construction of the LHC because the

- detector performance was

underestimated

MOTTO

*Rough estimates of BSM effects
can easily be deceptive*

example: discovery of the Higgs particle came much faster than expected at the time of construction of the LHC because the

- detector performance was
- theoretical prediction for Higgs production was significantly

underestimated

Status of particle physics: energy frontier

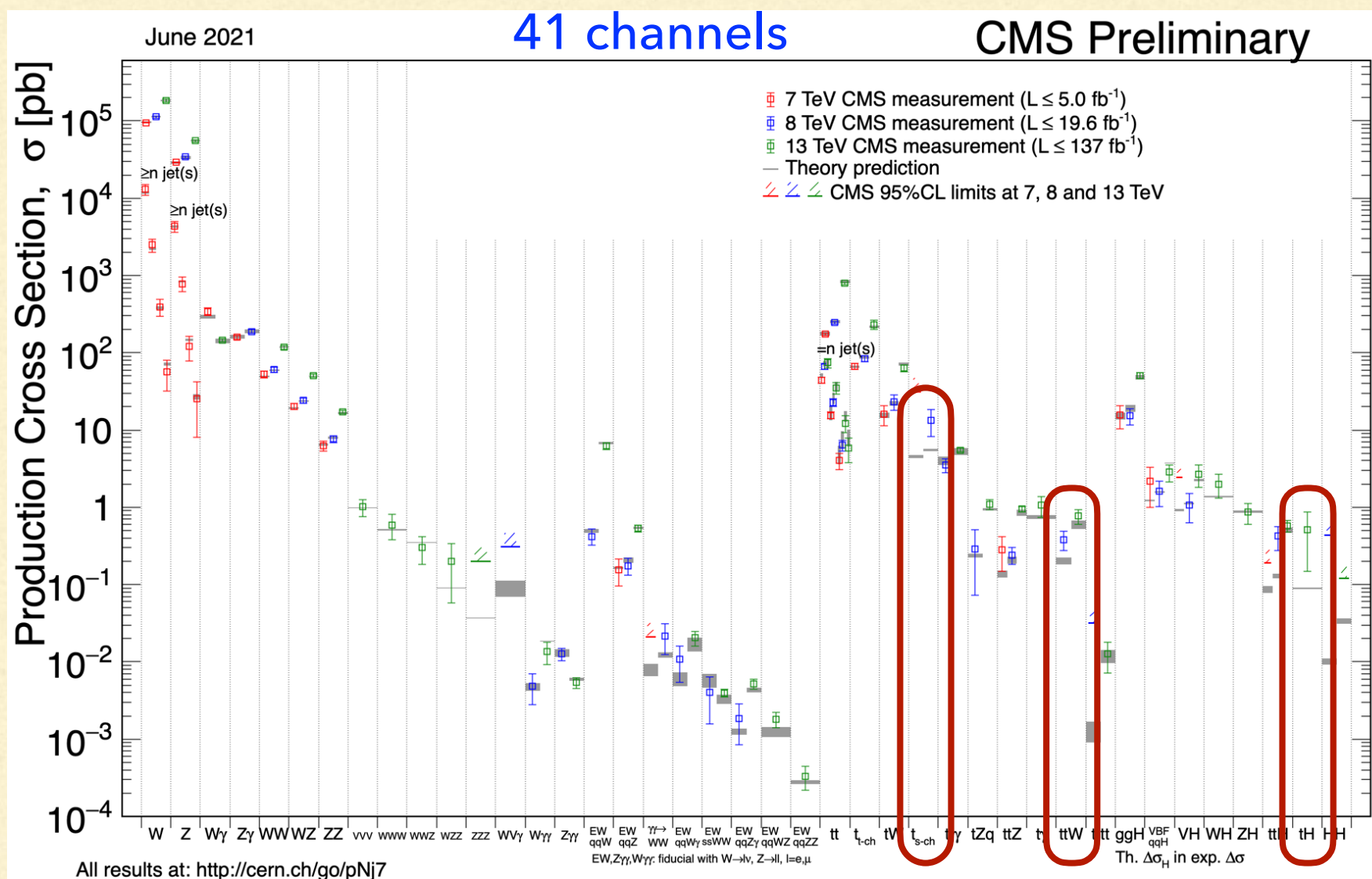
- Colliders: SM describes final states of particle collisions precisely
- No proven sign of new physics beyond SM at colliders*

*Exciting news keep popping up, all below discovery significance yet

Status of particle physics: energy frontier

- Colliders: SM describes final states of particle collisions precisely [CM]

[CMS public]



Status of particle physics: energy frontier

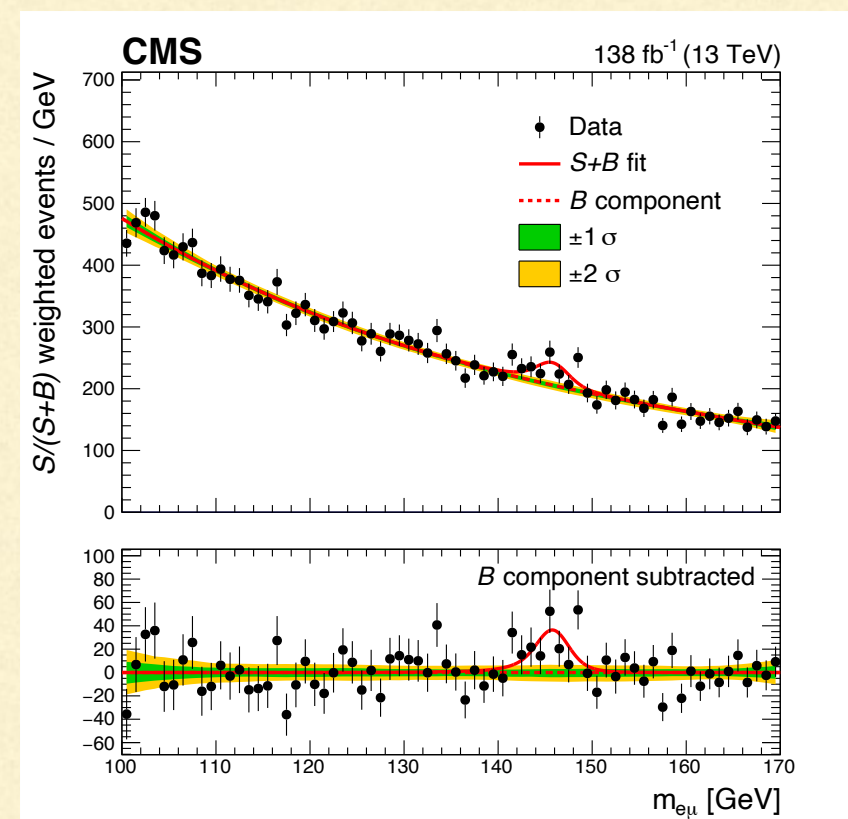
- Colliders: SM describes final states of particle collisions precisely
- No proven sign of new physics beyond SM at colliders

Status of particle physics: energy frontier

- Colliders: SM describes final states of particle collisions precisely
- No proven sign of new physics beyond SM at colliders*

$$pp \rightarrow X(= \text{new Higgs boson}) \rightarrow e^{\pm} \mu^{\mp}$$

[CMS preprint]



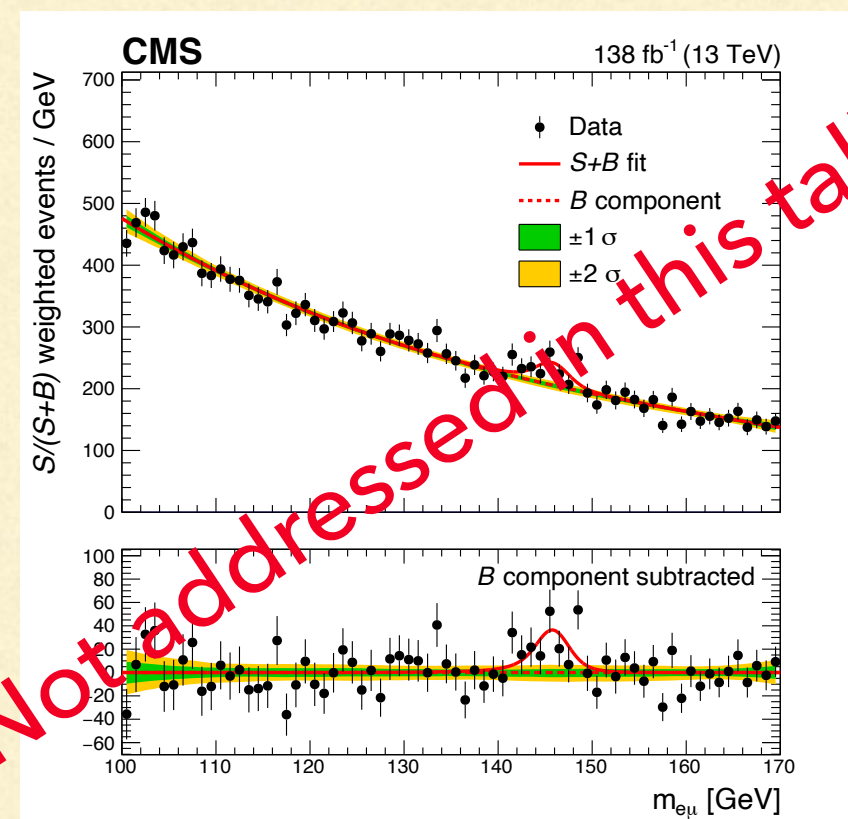
*Exciting news keep popping up, all below discovery significance yet

Status of particle physics: energy frontier

- Colliders: SM describes final states of particle collisions precisely
- No proven sign of new physics beyond SM at colliders*

$$pp \rightarrow X(= \text{new Higgs boson}) \rightarrow e^\pm \mu^\mp$$

[CMS preprint]



*Exciting news keep popping up, all below discovery significance yet

Status of particle physics: energy frontier

- Colliders: SM describes final states of particle collisions precisely
- No proven sign of new physics beyond SM at colliders
- SM vacuum is metastable
[Bezrukov et al, arXiv:1205.2893; Degrassi et al, arXiv:1205.6497]

Status of particle physics: cosmic and intensity frontiers

- Universe at large scale described precisely by cosmological SM:
 Λ CDM ($\Omega_m = 0.3$)
- Neutrino flavours oscillate
- Existing baryon asymmetry cannot be explained by CP asymmetry in SM
- Inflation of the early, accelerated expansion of the present Universe

[<https://pdg.lbl.gov>]

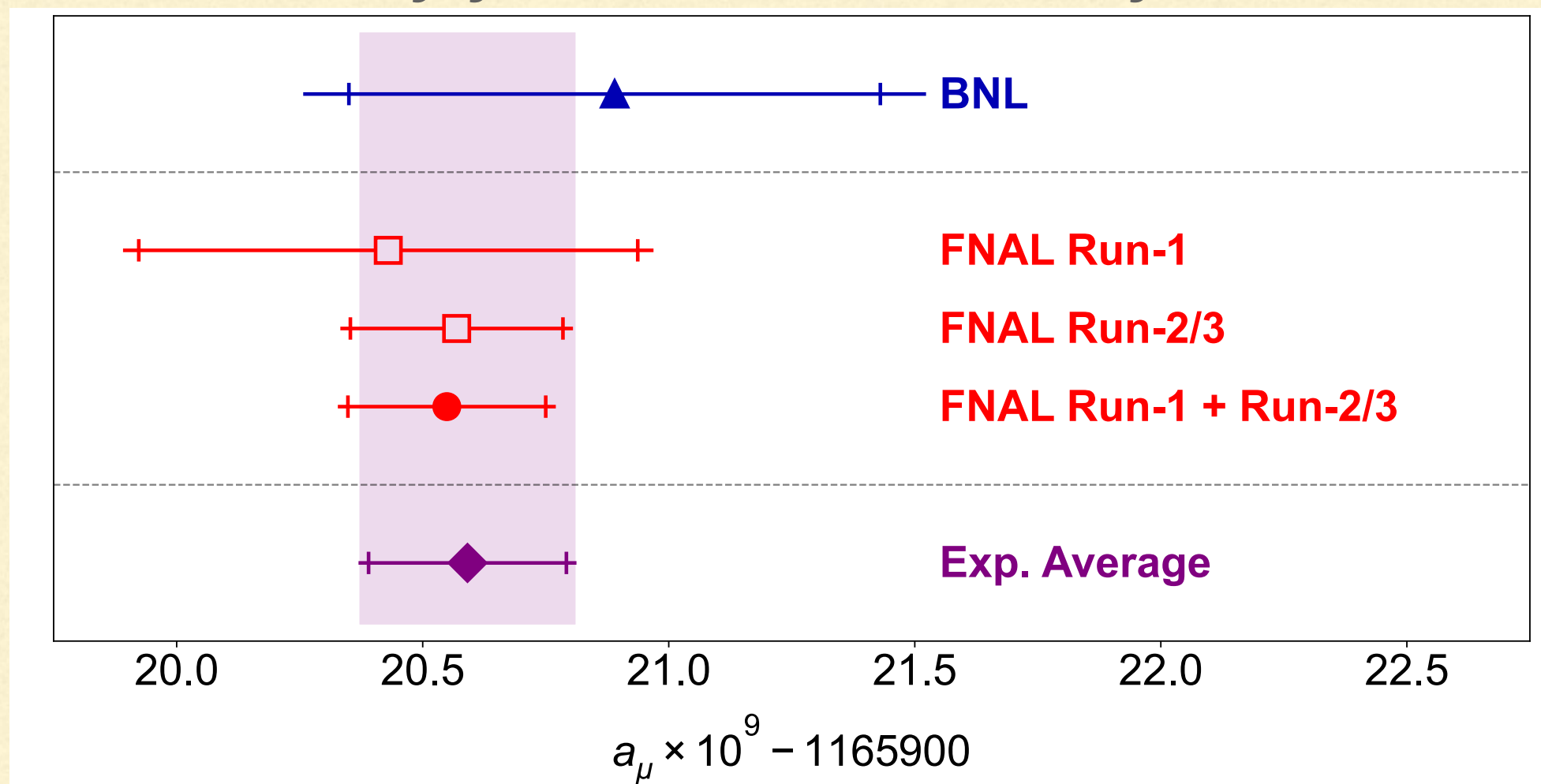
Established observations require physics beyond SM,
but do not suggest rich BSM physics

Phenomenological approach to new physics

Can we explain these observations,
but not more,
by the same model?

Status of the muon anomalous magnetic moment: experiment

- The muon g-2 has been a smoking gun for new physics for many years, more recently:



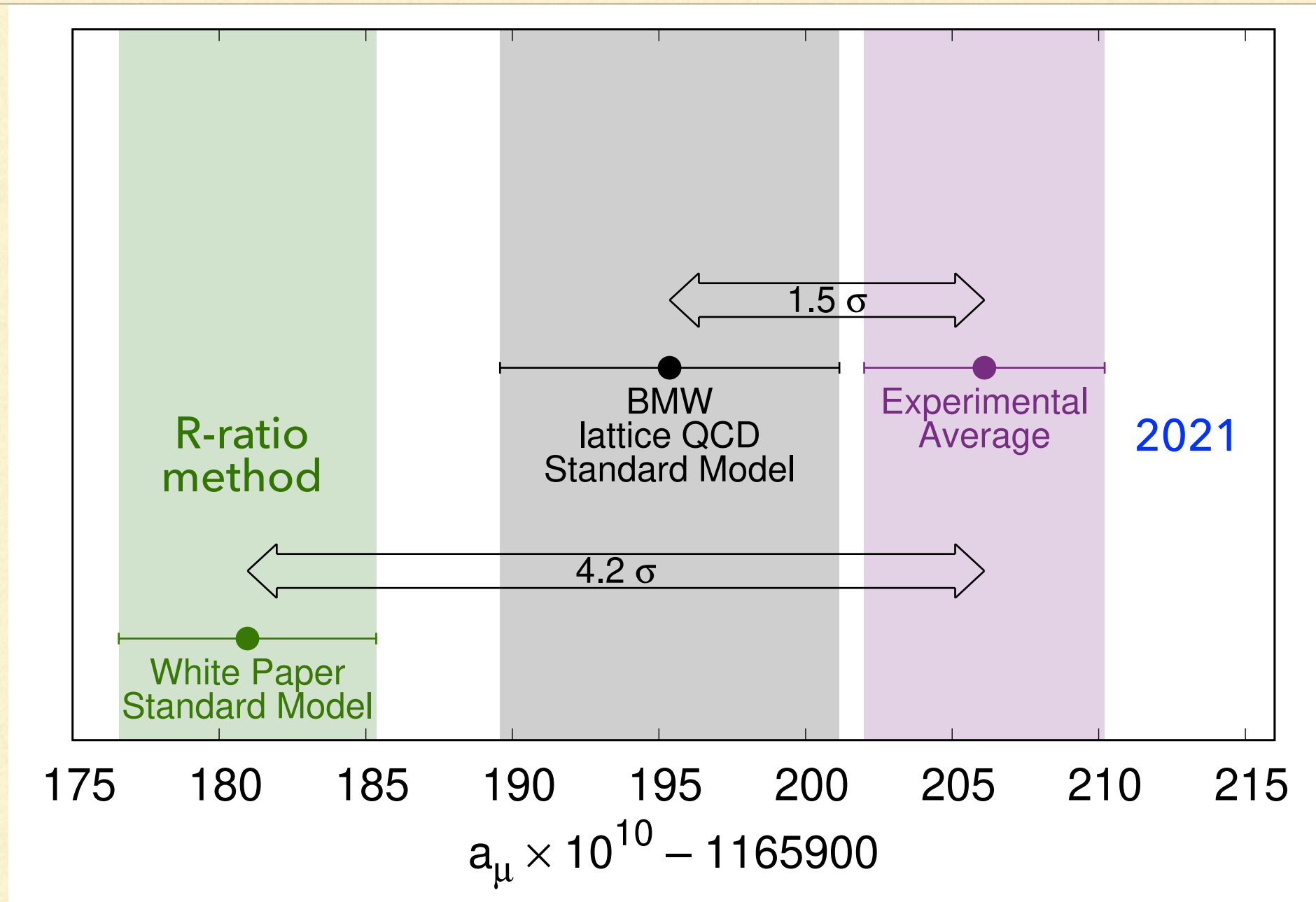
[<https://muon-g-2.fnal.gov/result2023.pdf>]

Status of the muon anomalous magnetic moment: experiment

- The muon g-2 has been a smoking gun for new physics for many years
- The most precise experimental value is from FNAL (2023) :
$$a_{\mu} = \frac{g - 2}{2} = 116\,592\,055(24) \cdot 10^{-11} \quad (0.20 \text{ ppm})$$
- ...equivalent to a bathroom scale sensitive to a single eyelash:

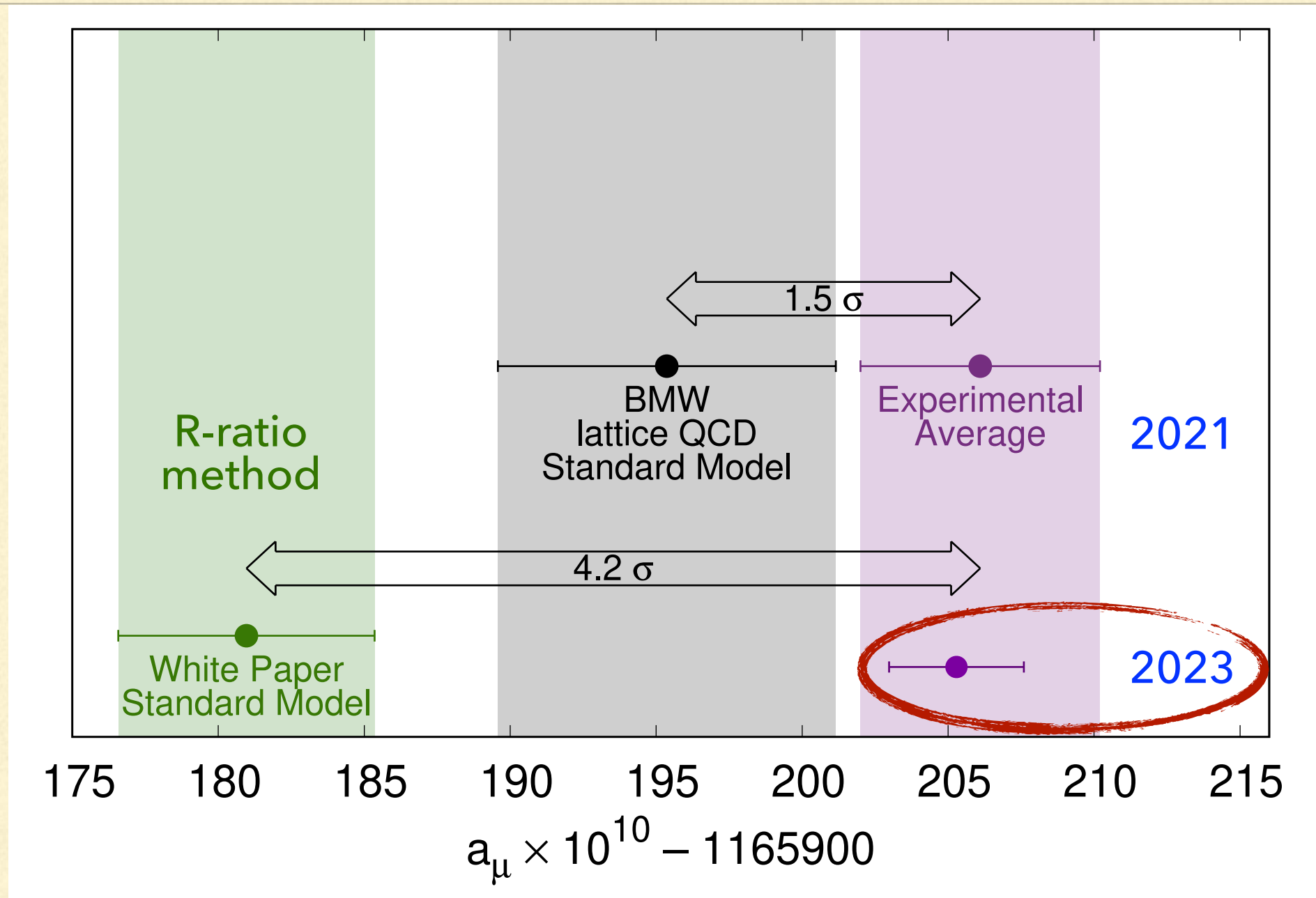


Status of the muon anomalous magnetic moment: experiment vs. theory



[BMW compilation]

Status of the muon anomalous magnetic moment: experiment vs. theory

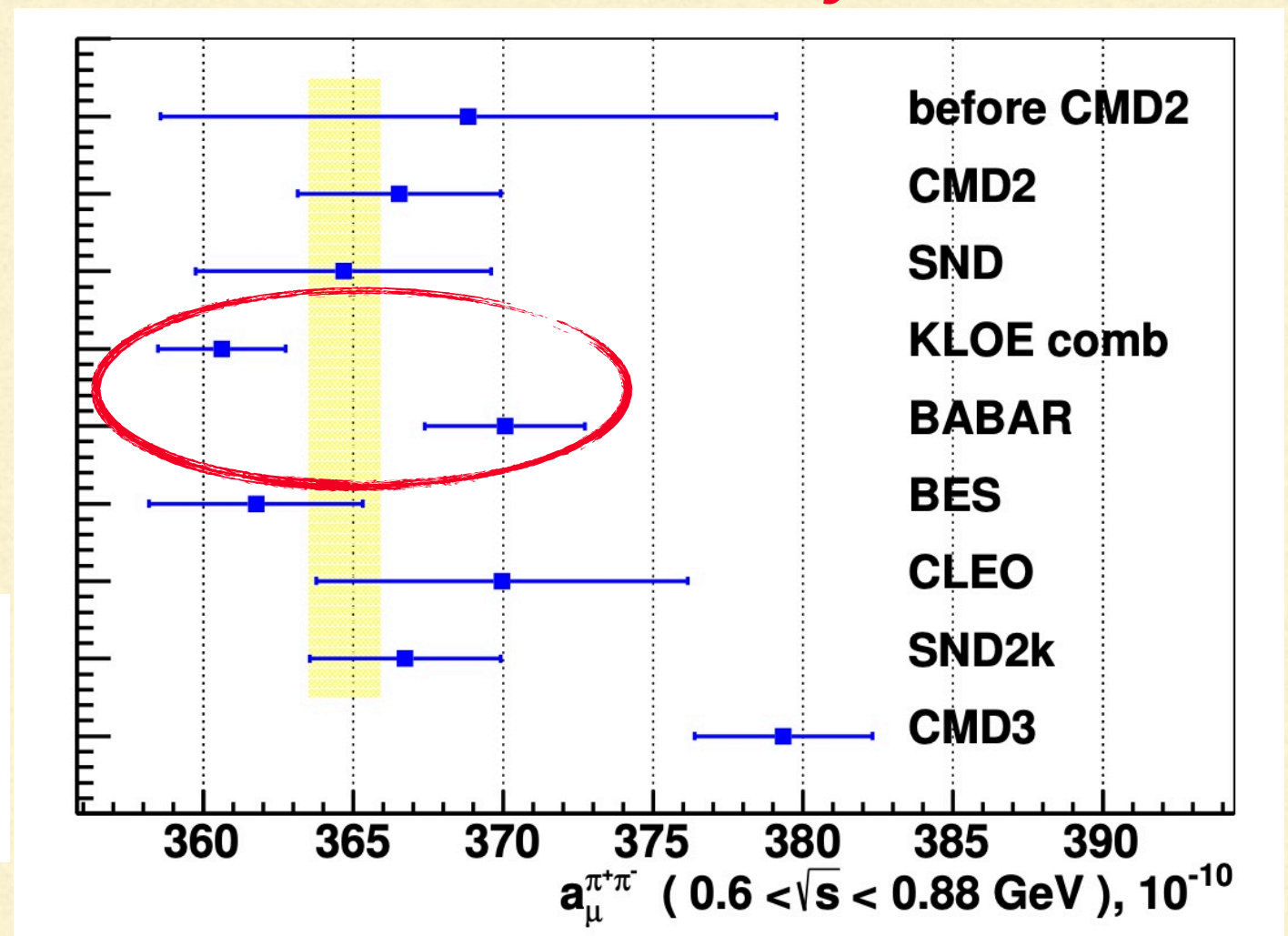
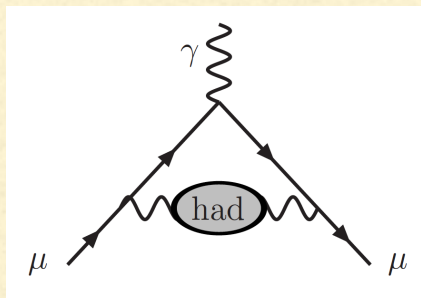


[BMW compilation]

Status of the muon anomalous magnetic moment: theory with R-ratio

- The muon g-2 has been a smoking gun for new physics for many years, **but tension already in earlier data used for theory prediction:**

$\sigma(e^+e^- \rightarrow \pi^+\pi^-)$ cross section in this energy range gives **more than 50%** to total **HVP** contribution to a_μ

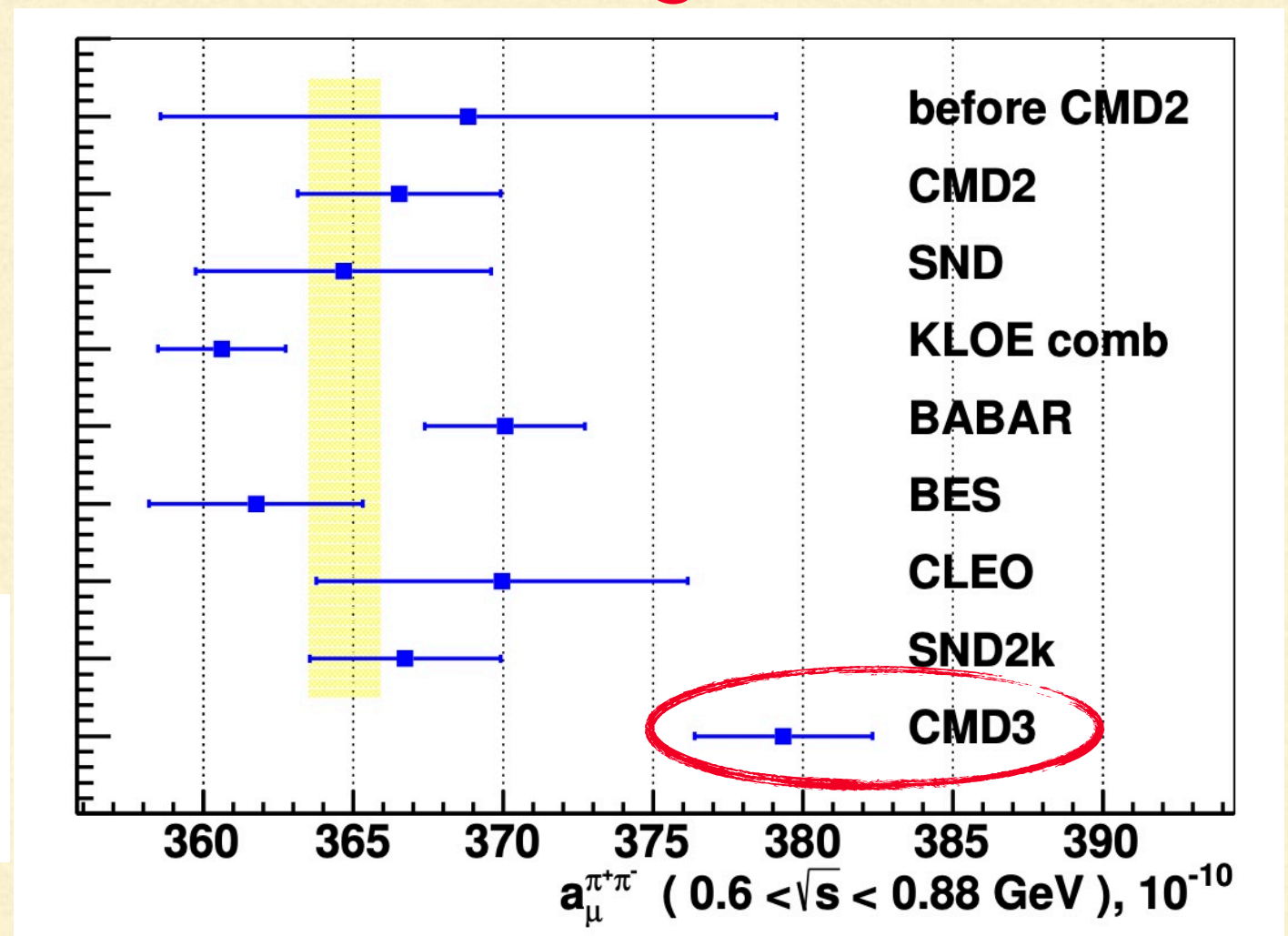
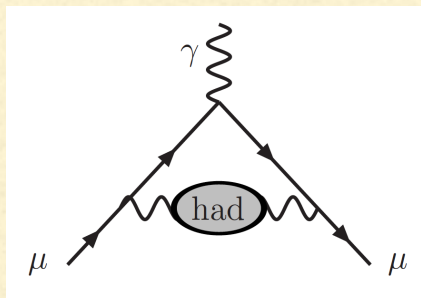


[BMW compilation]₁₇

Status of the muon anomalous magnetic moment: theory with R-ratio

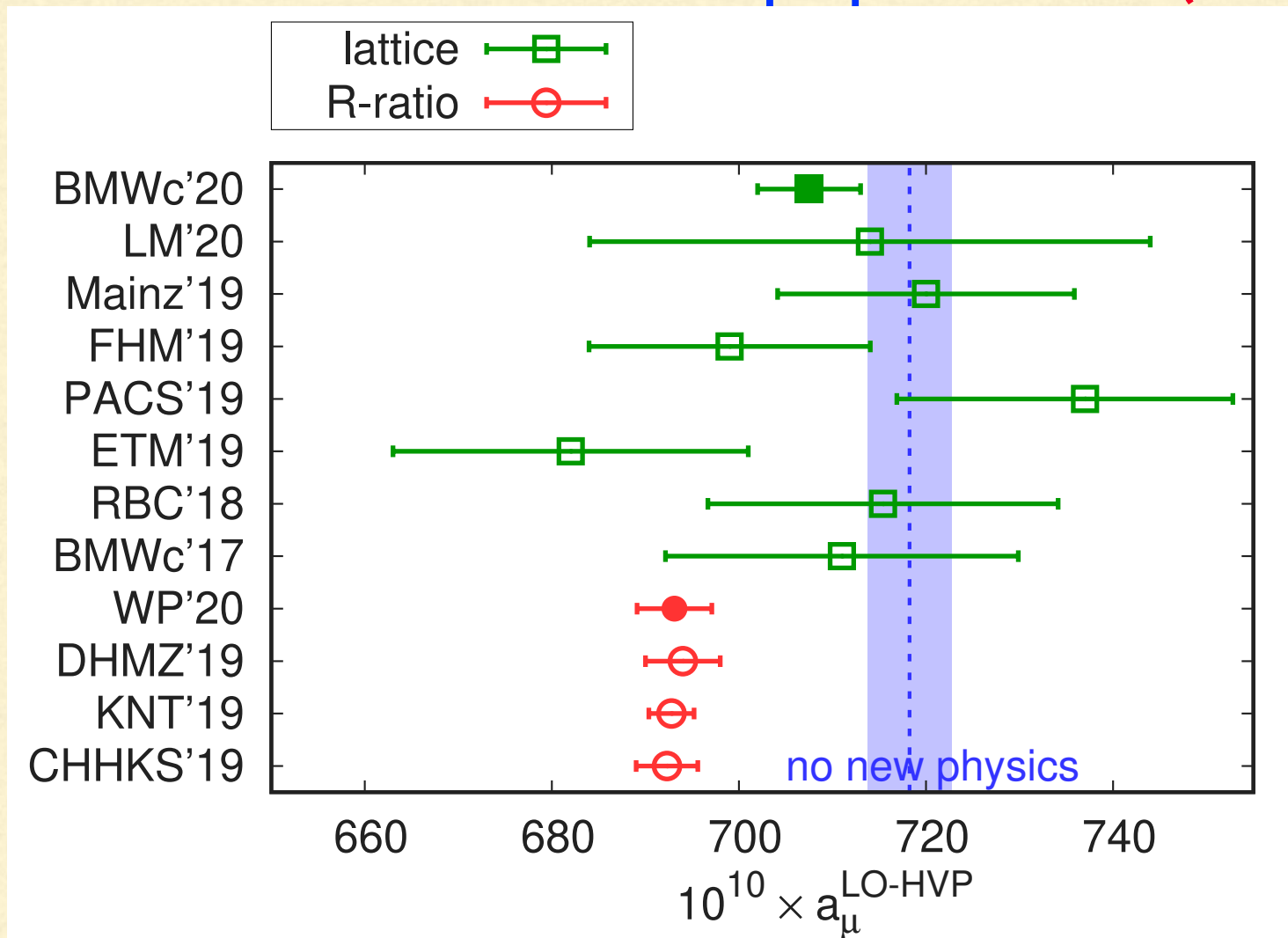
- New CMD3 data show a **~15 unit increase** in central value and **4.4 σ tension** with old average:

$\sigma(e^+e^- \rightarrow \pi^+\pi^-)$ cross section in this energy range gives **more than 50%** to total **HVP** contribution to a_μ



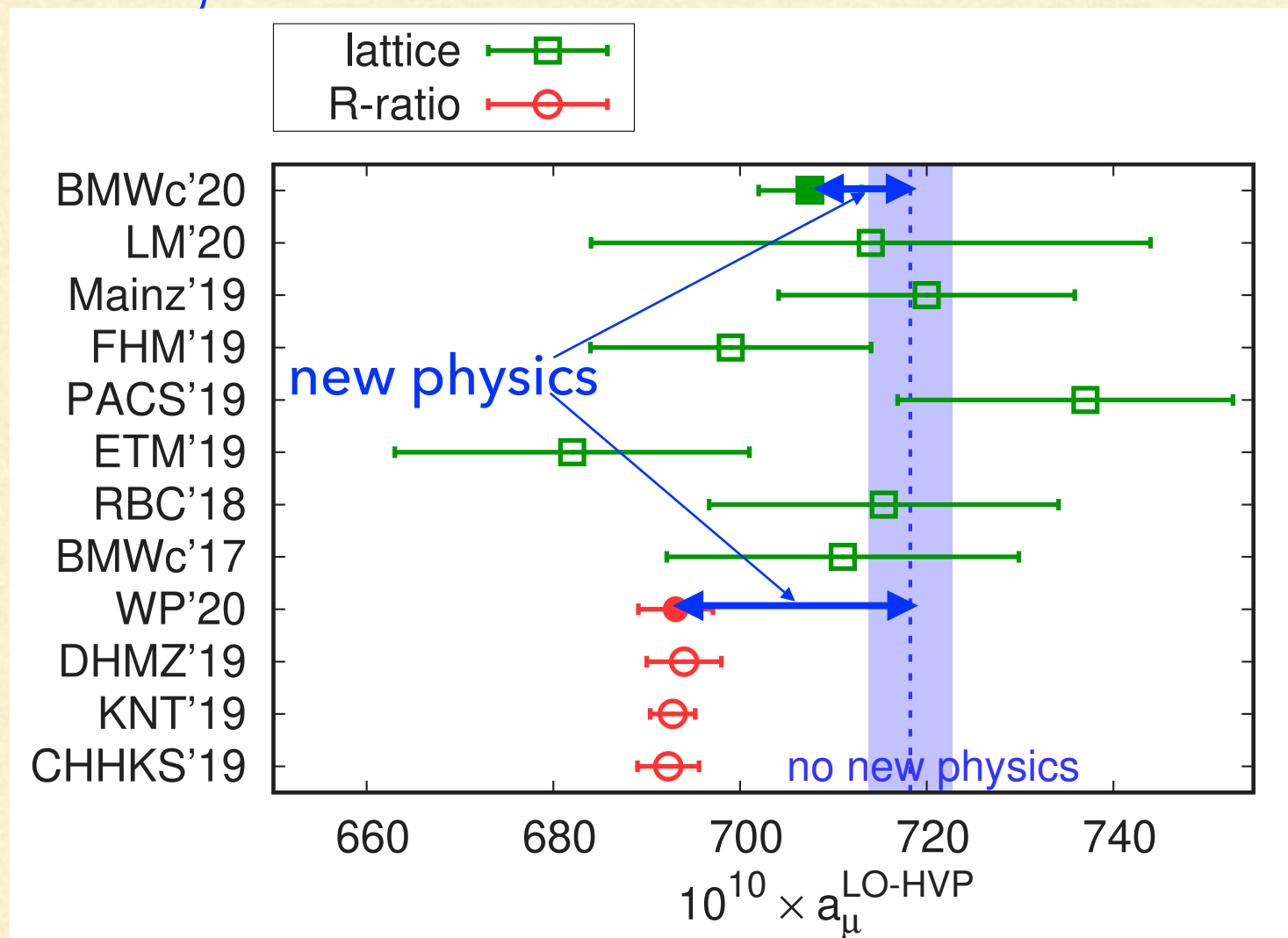
Status of the muon anomalous magnetic moment: lattice vs. R-ratio

- Lattice: $a_\mu^{\text{HVP}@LO} = 707.5(2.3)_{\text{stat}}(5.0)_{\text{sys}}[5.5]_{\text{tot}}$
- ~15 units above the R-ratio white paper value (a 2.1σ tension)



Message of the muon anomalous magnetic moment

- We are certain that there is new physics beyond the SM
- “Final word” on a_μ will tell how BSM should affect the muon g-2



Status of the muon anomalous magnetic moment

- We are certain that there is new physics beyond the SM
- Current main question:
 - How large is the new physics contribution to a_μ really?
 - "large" (almost 5σ – R-ratio result)
 - "small" (almost insignificant – lattice result)
- The experimental result appears robust, only its uncertainty will reduce further
- Main task:
 - Resolve discrepancy between theory predictions
- Until then
 - everything else is speculation

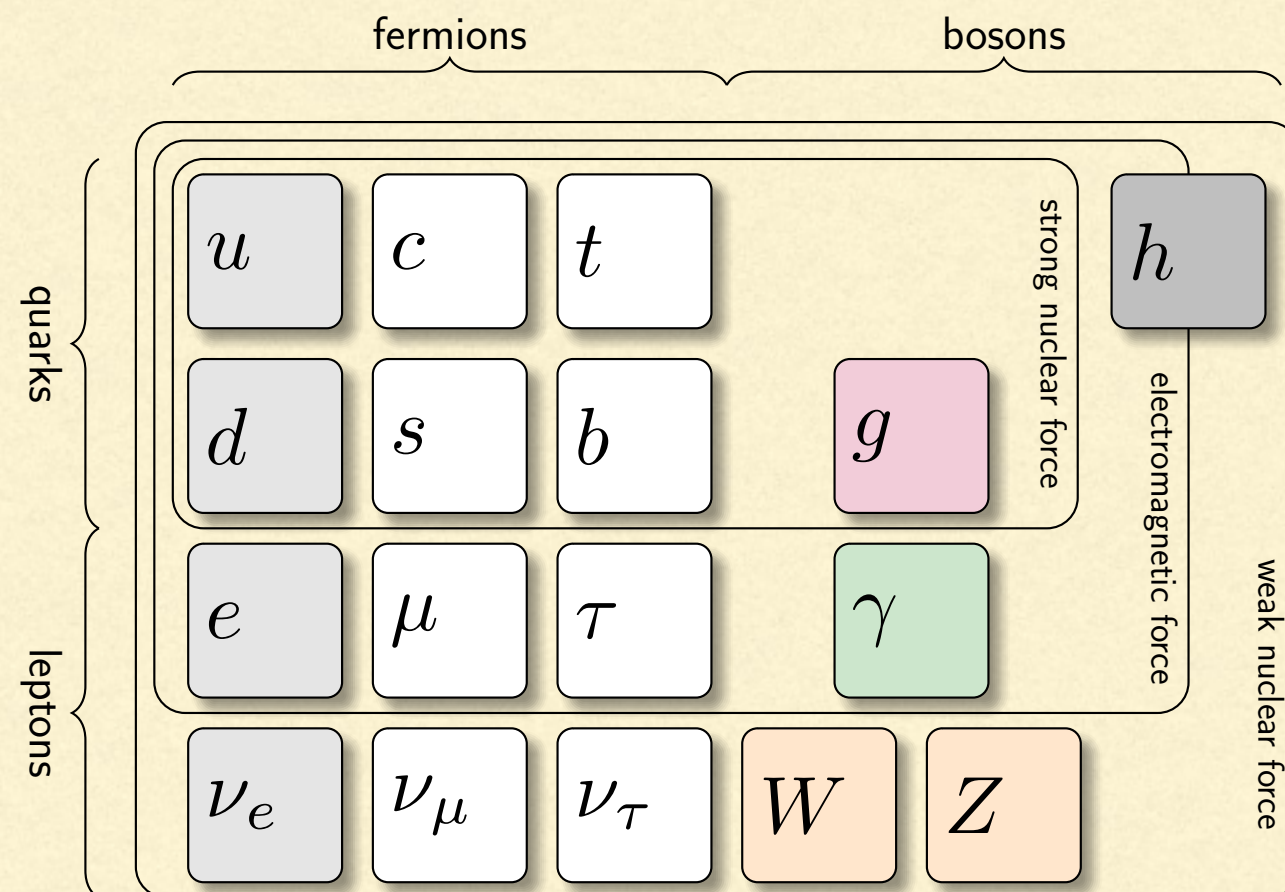
Muon anomalous magnetic moment: complying with lattice result

- New physics should have a small (smaller than EW) contribution to a_μ
- May constrain the available parameter space, but unlikely to exclude a model compatible with ElectroWeak Precision Observables (EWPOs)

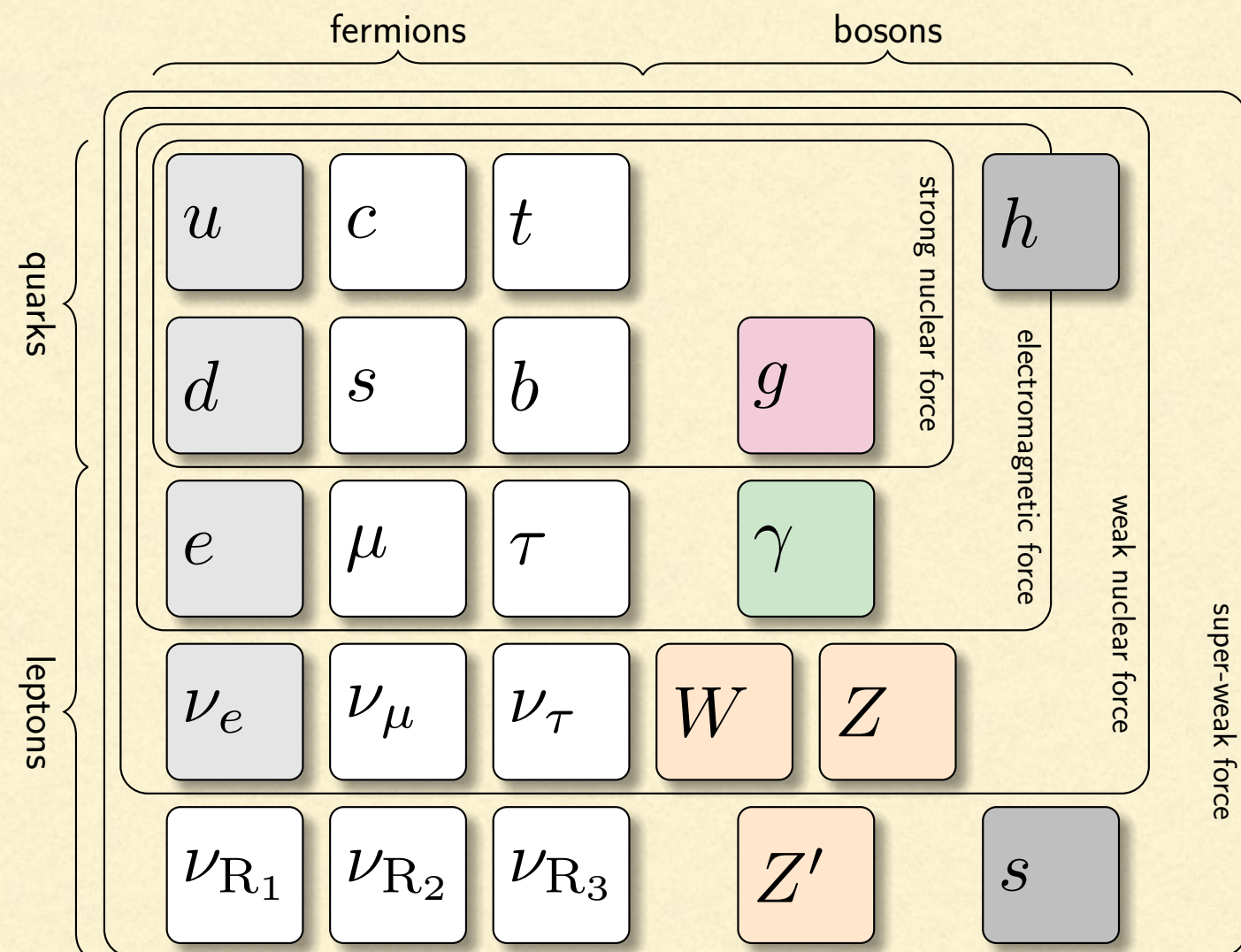
Extension of SM: three alternatives with different **strength** and **weaknesses**

- Effective field theory, such as **SMEFT**: **general** but **highly complex** (**2499** dim 6 operators), **focuses on new physics at high scales**
- Simplified models, such as **dark photon**, **extended scalar sector** or **right-handed neutrinos**: “**easily accessible**” **phenomenology**, but focus on specific aspect of new physics, so **cannot explain all BSM phenomena**
- UV complete extension with **potential of explaining BSM phenomena within a single model** such as **SuperWeak** extension of the **Standard Model**: **SWSM**

Particle content of SM



Particle content of SWSM (take-home picture)



Superweak extension of SM (SWSM)

- **Symmetry of the Lagrangian:** local
 $G = G_{\text{SM}} \times U(1)_Z$ with $G_{\text{SM}} = SU(3)_c \times SU(2)_L \times U(1)_Y$
renormalizable gauge theory, including all dim 4 operators allowed by G
- z-charges fixed by requirement of
 - gauge and gravity **anomaly cancellation** and
 - **gauge invariant Yukawa terms for neutrino mass generation**

Charge assignment from gauge invariant neutrino interactions

field	$SU(3)_c$	$SU(2)_L$	y_j	$z_j^{(a)}$	$z_j^{(b)}$	$r_j = z_j/z_\phi - y_j^{(c)}$
U_L, D_L	3	2	$\frac{1}{6}$	Z_1	$\frac{1}{6}$	0
U_R	3	1	$\frac{2}{3}$	Z_2	$\frac{7}{6}$	$\frac{1}{2}$
D_R	3	1	$-\frac{1}{3}$	$2Z_1 - Z_2$	$-\frac{5}{6}$	$-\frac{1}{2}$
ν_L, ℓ_L	1	2	$-\frac{1}{2}$	$-3Z_1$	$-\frac{1}{2}$	0
ν_R	1	1	0	$Z_2 - 4Z_1$	$\frac{1}{2}$	$\frac{1}{2}$
ℓ_R	1	1	-1	$-2Z_1 - Z_2$	$-\frac{3}{2}$	$-\frac{1}{2}$
ϕ	1	2	$\frac{1}{2}$	z_ϕ	1	$\frac{1}{2}$
χ	1	1	0	z_χ	-1	-1

Mixing in the neutral gauge sector

$$\begin{pmatrix} B_\mu \\ W_\mu^3 \\ B'_\mu \end{pmatrix} = \begin{pmatrix} c_W & -s_W & 0 \\ s_W & c_W & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_Z & -s_Z \\ 0 & s_Z & c_Z \end{pmatrix} \begin{pmatrix} A_\mu \\ Z_\mu \\ Z'_\mu \end{pmatrix} \quad \begin{aligned} c_X &= \cos \theta_X \\ s_X &= \sin \theta_X \end{aligned}$$

where θ_W is the weak mixing angle & θ_Z is the $Z - Z'$ mixing, implicitly:

$\tan(2\theta_Z) = -2\kappa / (1 - \kappa^2 - \tau^2)$, with κ and τ effective couplings,
functions of the Lagrangian couplings

The expressions for the neutral gauge boson masses are somewhat cumbersome, but exists a nice, **compact generalization** of the **SM**

mass-relation formula: $\frac{M_W^2}{c_W^2} = c_Z^2 M_Z^2 + s_Z^2 M_{Z'}^2$ $\left(M_W = \frac{1}{2} g_L v \right)$

Scalars in the SWSM

- Standard ϕ complex $SU(2)_L$ doublet and new χ complex singlet:

$$\mathcal{L}_{\phi,\chi} = [D_{\mu}^{(\phi)} \phi]^* D^{(\phi)\mu} \phi + [D_{\mu}^{(\chi)} \chi]^* D^{(\chi)\mu} \chi - V(\phi, \chi)$$

- with scalar potential

$$V(\phi, \chi) = V_0 - \mu_{\phi}^2 |\phi|^2 - \mu_{\chi}^2 |\chi|^2 + (|\phi|^2, |\chi|^2) \begin{pmatrix} \lambda_{\phi} & \frac{\lambda}{2} \\ \frac{\lambda}{2} & \lambda_{\chi} \end{pmatrix} \begin{pmatrix} |\phi|^2 \\ |\chi|^2 \end{pmatrix}$$

- After SSB, $G \rightarrow SU(3)_c \times U(1)_{QED}$ in R_{ξ} gauge

$$\phi = \frac{1}{\sqrt{2}} \begin{pmatrix} -i\sqrt{2}\sigma^+ \\ v + h' + i\sigma_{\phi} \end{pmatrix} \quad \& \quad \chi = \frac{1}{\sqrt{2}} (w + s' + i\sigma_{\chi})$$

Mixing in the scalar sector

$$\begin{pmatrix} h' \\ s' \end{pmatrix} = \begin{pmatrix} c_S & s_S \\ -s_S & c_S \end{pmatrix} \begin{pmatrix} h \\ s \end{pmatrix}$$

where θ_S is the scalar mixing angle implicitly:

$$\tan(2\theta_S) = \lambda v w / \left(\lambda_\chi w^2 - \lambda_\phi v^2 \right), \text{ with } v \text{ and } w \text{ VEVs}$$

5 new parameters:

- in **gauge** sector: $\{g_z \text{ and } g_{yz}\}$ or $\{\kappa \text{ and } \tau\}$ or $\{\theta_Z \text{ and } M_Z\}$
- in **scalar** sector: $\{\mu_\chi^2, \lambda_\chi \text{ and } \lambda\}$ or $\{w, \lambda_\chi \text{ and } \lambda\}$ or $\{M_S, \theta_S \text{ and } \lambda\}$

After SSB neutrino mass terms appear

$$-\mathcal{L}_Y^\ell = \frac{w + s' + i\sigma_\chi}{2\sqrt{2}} \overline{\nu_R^c} \mathbf{Y}_N \nu_R + \frac{v + h' - i\sigma_\phi}{\sqrt{2}} \overline{\nu_L} \mathbf{Y}_\nu \nu_R + \text{h.c.}$$

$$\mathbf{M}_N = \frac{w}{\sqrt{2}} \mathbf{Y}_N$$

$$\mathbf{M}_D = \frac{v}{\sqrt{2}} \mathbf{Y}_\nu$$

- In flavour basis the full 6×6 mass matrix reads $\mathbf{M}' = \begin{pmatrix} \mathbf{0}_3 & \mathbf{M}_D^T \\ \mathbf{M}_D & \mathbf{M}_N \end{pmatrix}$
- ν_L and ν_R have the same q-numbers, can mix, leading to type-I see-saw
- Dirac and Majorana mass terms appear already at tree level by SSB (not generated radiatively)
- Quantum corrections to active neutrinos are not dangerous
[Iwamoto et al, [arXiv:2104.14571](https://arxiv.org/abs/2104.14571)]

Expected consequences (take-home messages)

- Dirac and Majorana neutrino mass terms are generated by the SSB of the scalar fields, providing the origin of neutrino masses and oscillations
[Iwamoto, Kärkkäinen, Péli, ZT, arXiv:[2104.14571](#); Kärkkäinen and ZT, arXiv:[2105.13360](#)]
- The lightest new particle is a natural and viable candidate for WIMP dark matter if it is sufficiently stable [Seller, Iwamoto and ZT, arXiv:[2104.11248](#)]
- Diagonalization of neutrino mass terms leads to the PMNS matrix, which in turn can be the source of lepto-baryogenesis
[Seller, Szép, ZT, arXiv:[2301.07961](#) and under investigation]
- The second scalar together with the established BEH field can stabilize the vacuum and be related to the accelerated expansion now and inflation in the early universe
[Péli, Nándori and ZT, arXiv:[1911.07082](#); Péli and ZT, arXiv:[2204.07100](#)]

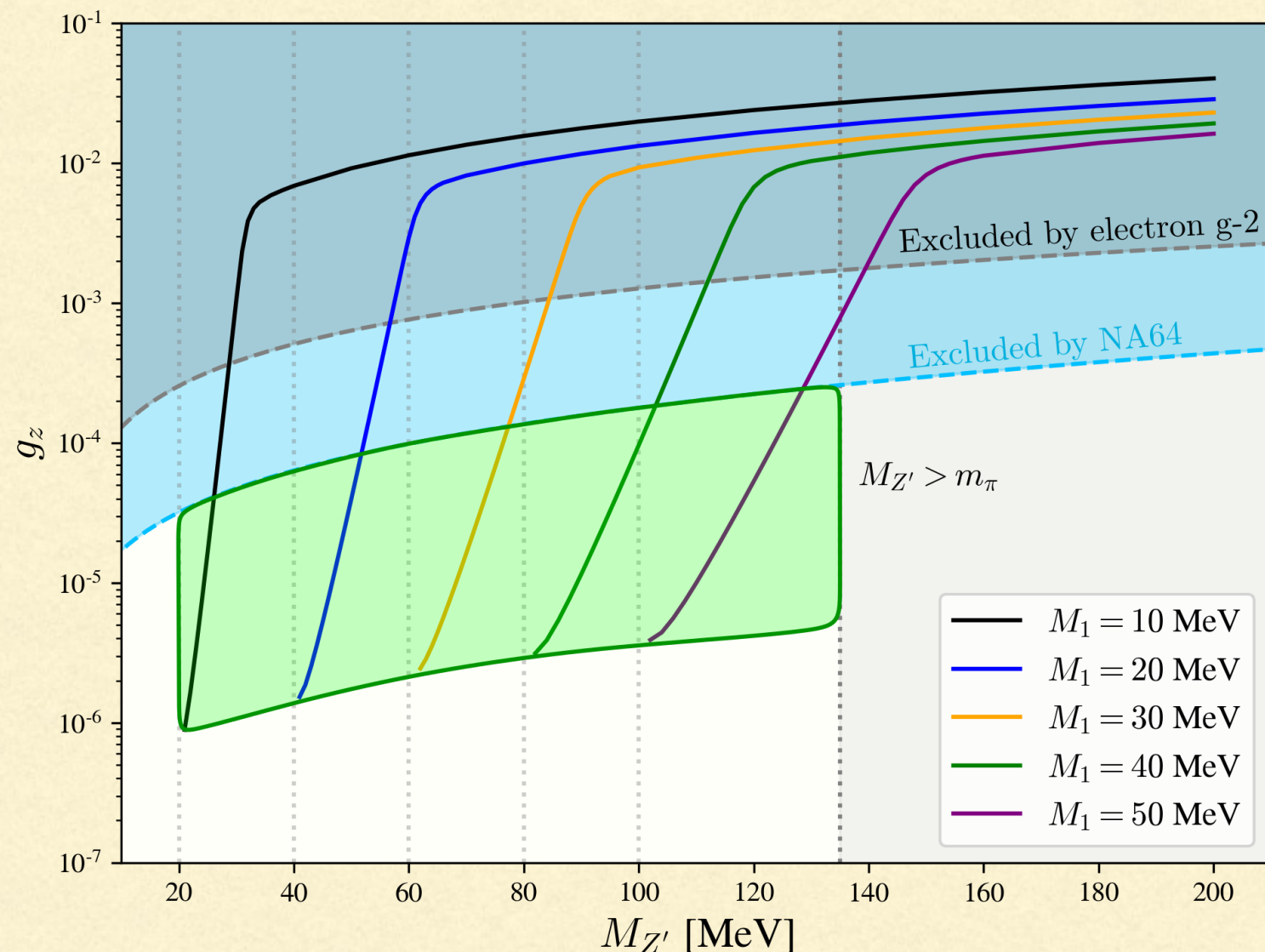
Expected consequences (take-home messages)

- Dirac and Majorana neutrino mass terms are generated by the SSB of the scalar fields, providing the origin of neutrino masses and oscillations
[Iwamoto, Kärkäinen, Péli, ZT, arXiv:[2104.14571](#); Kärkkäinen and ZT, arXiv:[2105.13360](#)]
- The lightest new particle is a natural and viable candidate for WIMP dark matter if it is sufficiently stable [Seller, Iwamoto and ZT, arXiv:[2104.11248](#)]
- Diagonalization of neutrino mass terms leads to the PMNS matrix, which in turn can be the source of lepto-baryogenesis
[Seller, Szép, ZT, arXiv:[2301.07961](#) and under investigation]
- The second scalar together with the established BEH field can stabilize the vacuum and be related to the accelerated expansion now and inflation in the early universe
[Péli, Nándori and ZT, arXiv:[1911.07082](#); Péli and ZT, arXiv:[2204.07100](#)]

Dark matter candidate

- **DM exists**, but known evidence is based solely on the gravitational effect of the dark matter on the luminous astronomical objects and on the Hubble-expansion of the Universe
- Assume that the DM has particle origin
- Only chance to observe such a particle if it **interacts with the SM particles, which needs a portal**
In the superweak model the vector boson portal Z' with the lightest sterile neutrino ν_4 as dark matter candidate is a natural scenario (Higgs portal exists, but negligible)

Parameter space for the freeze-out scenario of dark matter production in the SWSM

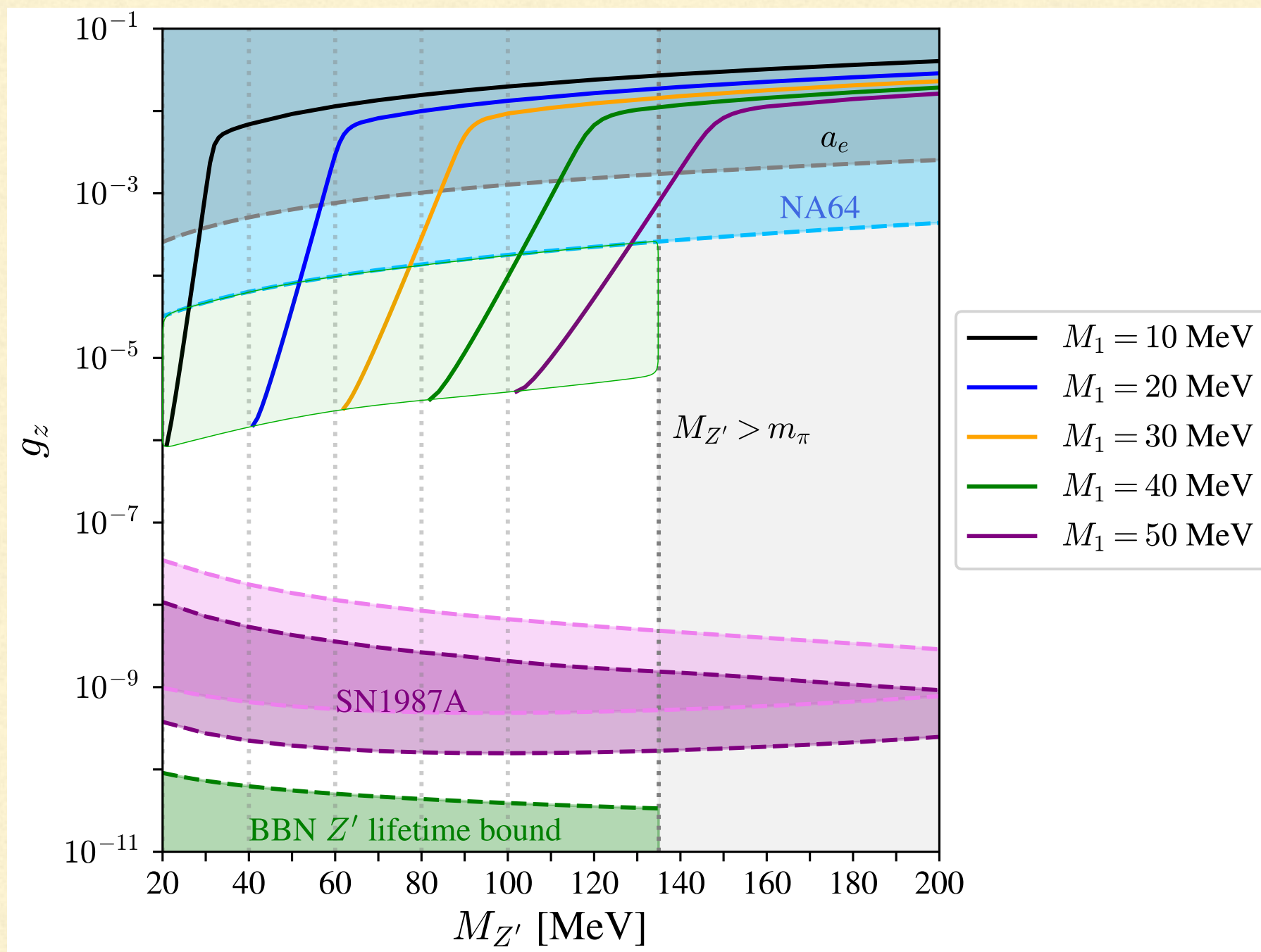


It is essential for the SWSM DM candidate that the resonance in $\text{SM}+\text{SM} \rightarrow Z' \rightarrow \text{DM}+\text{DM}$ can dominate the integral in the rate

Experimental constraints

- **Anomalous magnetic moment** of electron and muon
 - Z' couples to leptons modifying the magnetic moment
 - Constraints on $(g - 2)$ translate to upper bounds on the coupling $g_z(M_{Z'})$
- **NA64 search for missing energy events**
 - **Strict upper bounds** on $g_z(M_{Z'})$ for any U(1) extension (dark photons)
- **Supernova constraints** based on SN1987A
 - Constraints are based on comparing observed and calculated neutrino fluxes
- **Big Bang Nucleosynthesis** provides **constraints on new particles**
 - New particles should have negligible effects during BBN
 - Meson production can be dangerous close to BBN
- Further constraints are due to **CMB, solar cooling, beam dump experiments** etc.

Cosmological constraints on the freeze-out scenario of dark matter production in the SWSM

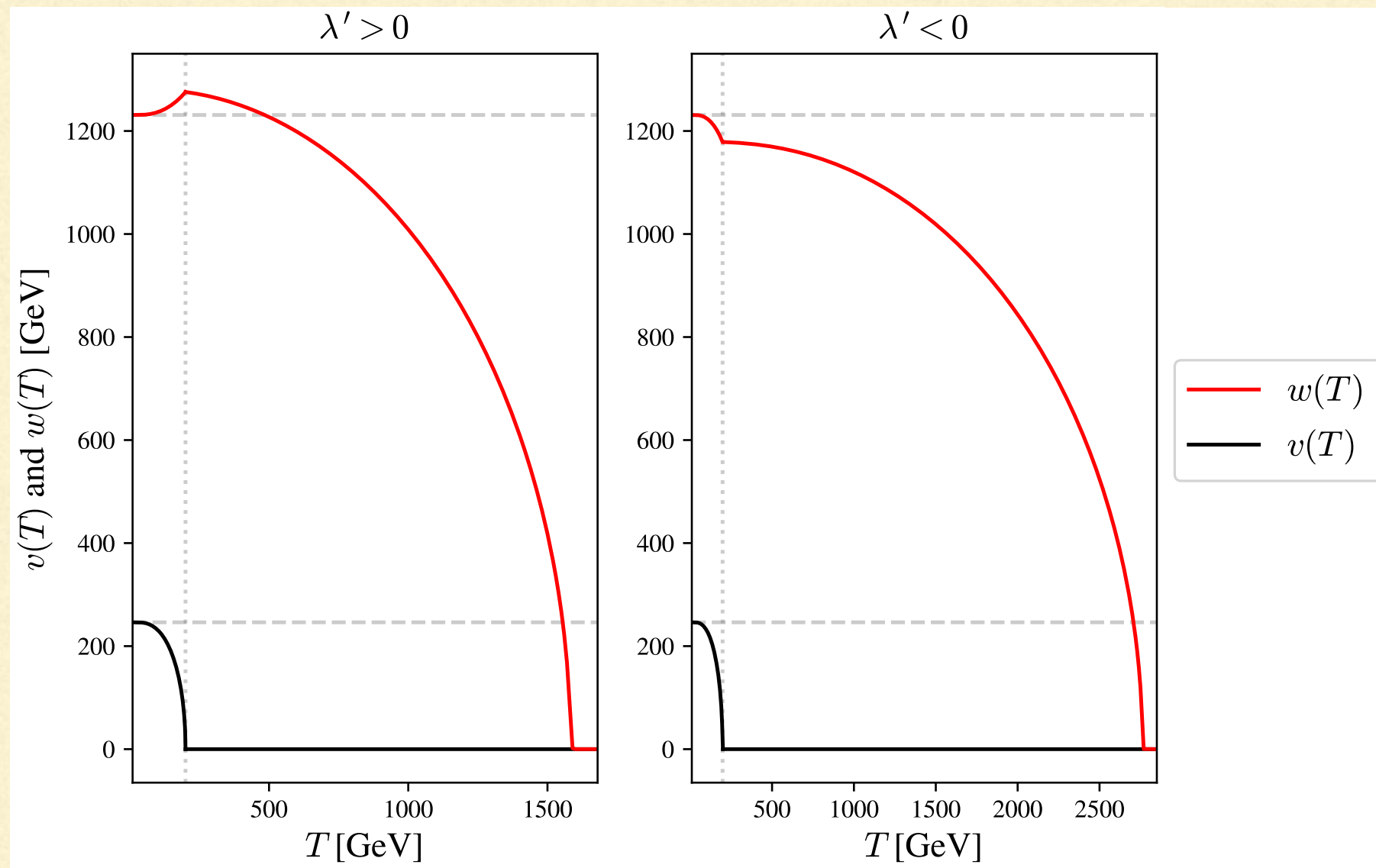


Expected consequences (take-home messages)

- Dirac and Majorana neutrino mass terms are generated by the SSB of the scalar fields, providing the origin of neutrino masses and oscillations
[Iwamoto, Kärkkäinen, Péli, ZT, arXiv:[2104.14571](#); Kärkkäinen and ZT, arXiv:[2105.13360](#)]
- The lightest new particle is a natural and viable candidate for WIMP dark matter if it is sufficiently stable [Seller, Iwamoto and ZT, arXiv:[2104.11248](#)]
- Diagonalization of neutrino mass terms leads to the PMNS matrix, which in turn can be the source of leptogenesis [Seller, Szép, ZT, arXiv:[2301.07961](#) and under investigation]
- The second scalar together with the established BEH field can stabilize the vacuum and be related to the accelerated expansion now and inflation in the early universe
[Péli, Nándori and ZT, arXiv:[1911.07082](#); Péli and ZT, arXiv:[2204.07100](#)]

Prerequisite: Phase-transitions in the SWSM

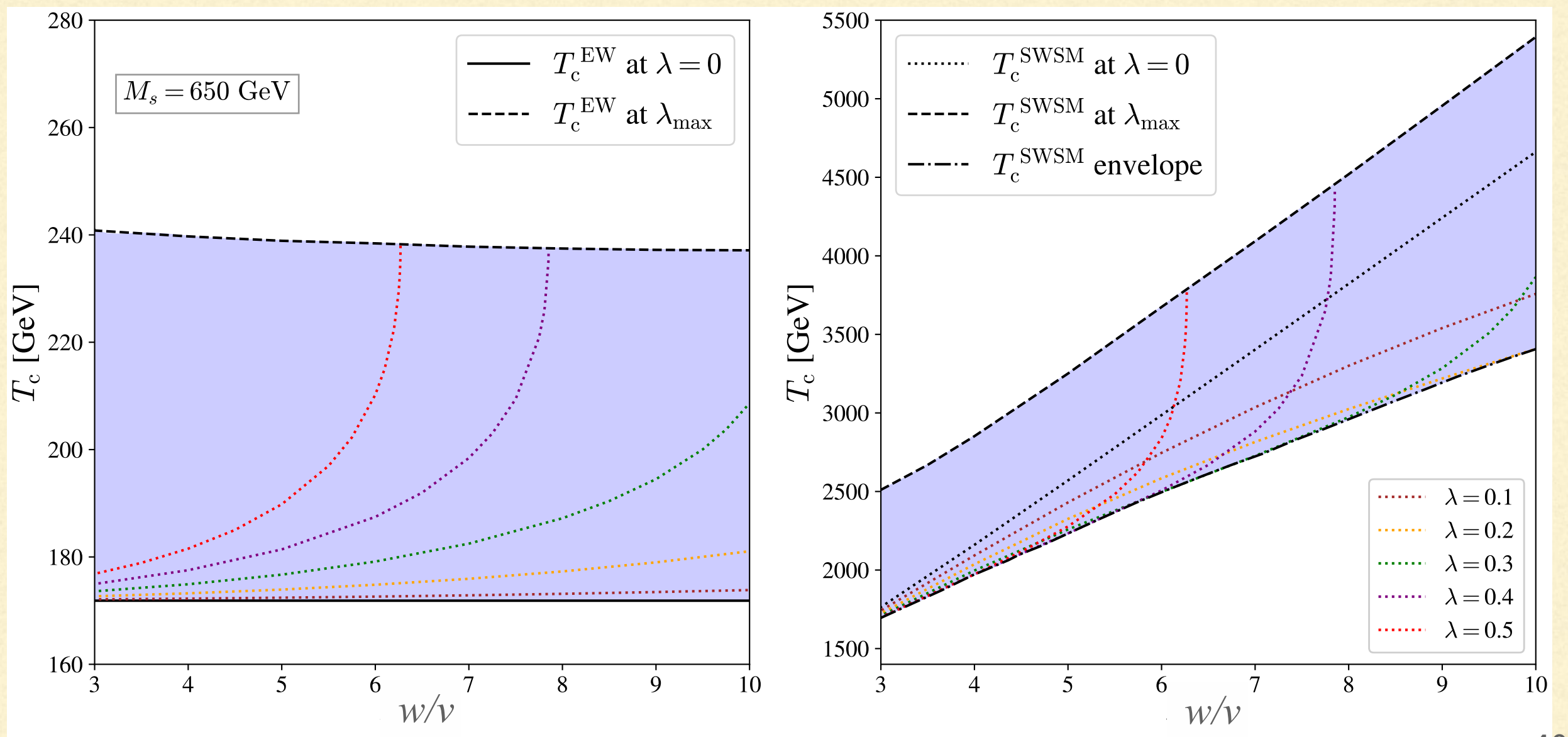
$U(1)_Z$ is broken earlier than $SU(2)_L \times U(1)_Y$



$$M_S = 200 \text{ GeV}, \quad M_N = 150 \text{ GeV}, \quad w = 5v, \quad |\lambda| = 0.0394$$

Prerequisite: phase-transition temperatures in the SWSM

$U(1)_Z$ is broken earlier than $SU(2)_L \times U(1)_Y$



Expected consequences (take-home messages)

- Dirac and Majorana neutrino mass terms are generated by the SSB of the scalar fields, providing the origin of neutrino masses and oscillations
[Iwamoto, Kärkäinen, Péli, ZT, [arXiv:2104.14571](#); Kärkkäinen and ZT, [arXiv:2105.13360](#)]
- The lightest new particle is a natural and viable candidate for WIMP dark matter if it is sufficiently stable
[Seller, Iwamoto and ZT, [arXiv:2104.11248](#)]
- Diagonalization of neutrino mass terms leads to the PMNS matrix, which in turn can be the source of lepto-baryogenesis
[Seller, Szép, ZT, [arXiv:2301.07961](#) and under investigation]
- The second scalar together with the established BEH field can stabilize the vacuum and be related to the accelerated expansion now and inflation in the early universe
[Péli, Nándori and ZT, [arXiv:1911.07082](#); Péli and ZT, [arXiv:2204.07100](#)]

Expected consequences (take-home messages)

- Dirac and Majorana neutrino mass terms are generated by the SSB of the scalar fields, providing the origin of neutrino masses and oscillations
[Iwamoto, Kärkkäinen, Péli, ZT, arXiv:[2104.14571](#); Kärkkäinen and ZT, arXiv:[2105.13360](#)]
- The lightest new particle is a natural and viable candidate for WIMP dark matter if it is sufficiently stable
[Seller, Iwamoto and ZT, arXiv:[2104.11248](#)]
- Diagonalization of neutrino mass terms leads to the PMNS matrix, which in turn can be the source of lepto-baryogenesis
[Seller, Szép, ZT, arXiv:[2301.07961](#) and under investigation]
- The second scalar together with the established BEH field can stabilize the vacuum and be related to the accelerated expansion now and inflation in the early universe
[Péli, Nándori and ZT, arXiv:[1911.07082](#); Péli and ZT, arXiv:[2204.07100](#)]

SWSM has the potential of explaining all known results beyond the SM

Main questions

Is there a non-empty region of the parameter space where all these promises are fulfilled?

Main questions

Is there a non-empty region of the parameter space where all these promises are fulfilled?

Can we predict any new phenomenon observable by present or future experiments?

Main questions

Present focus:

Is there a non-empty region of the parameter space where all these promises are fulfilled?

Can we predict any new phenomenon observable by present or future experiments?

Important test

Once the allowed region of the parameter space for fulfilling the expectations is understood

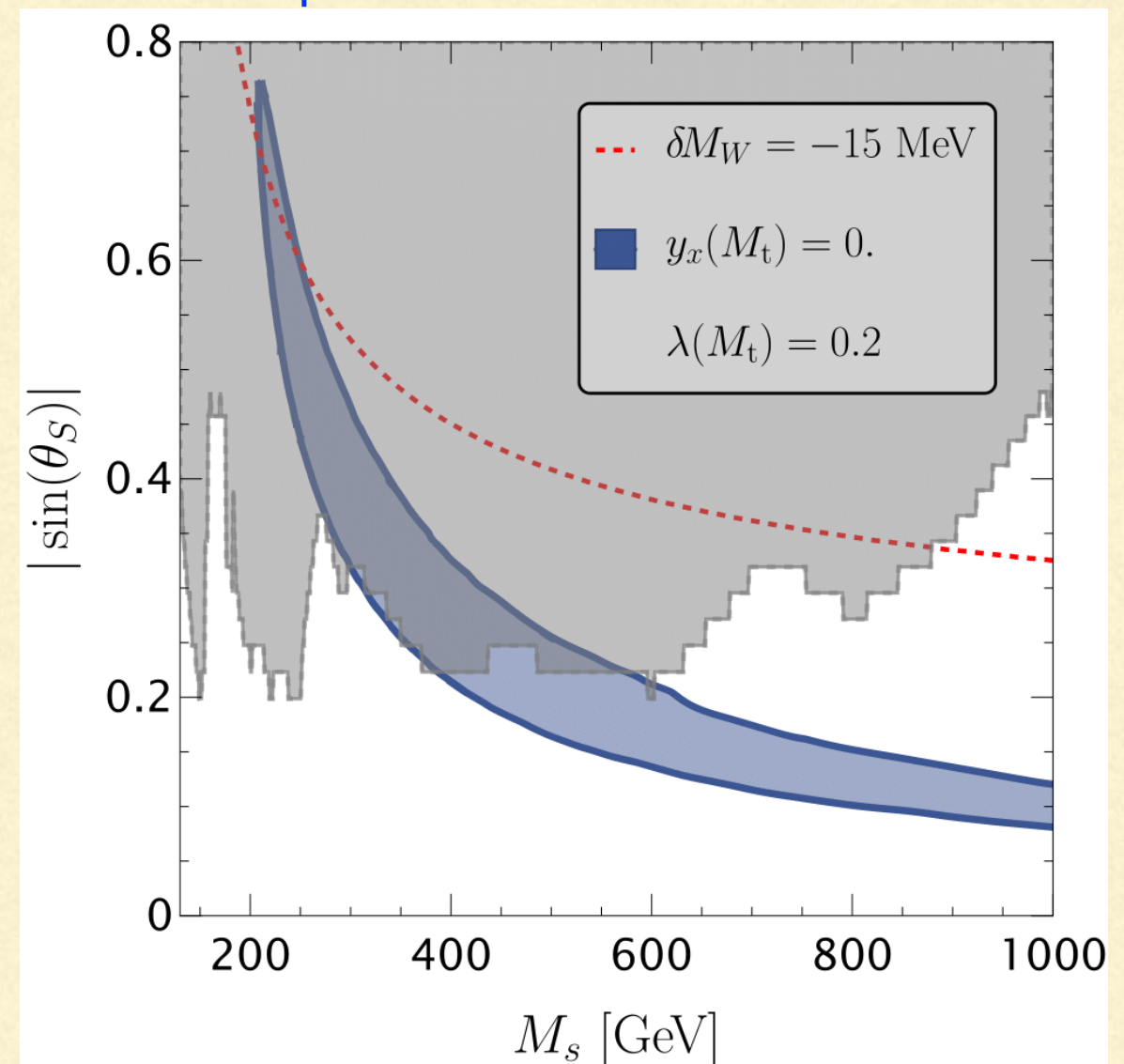
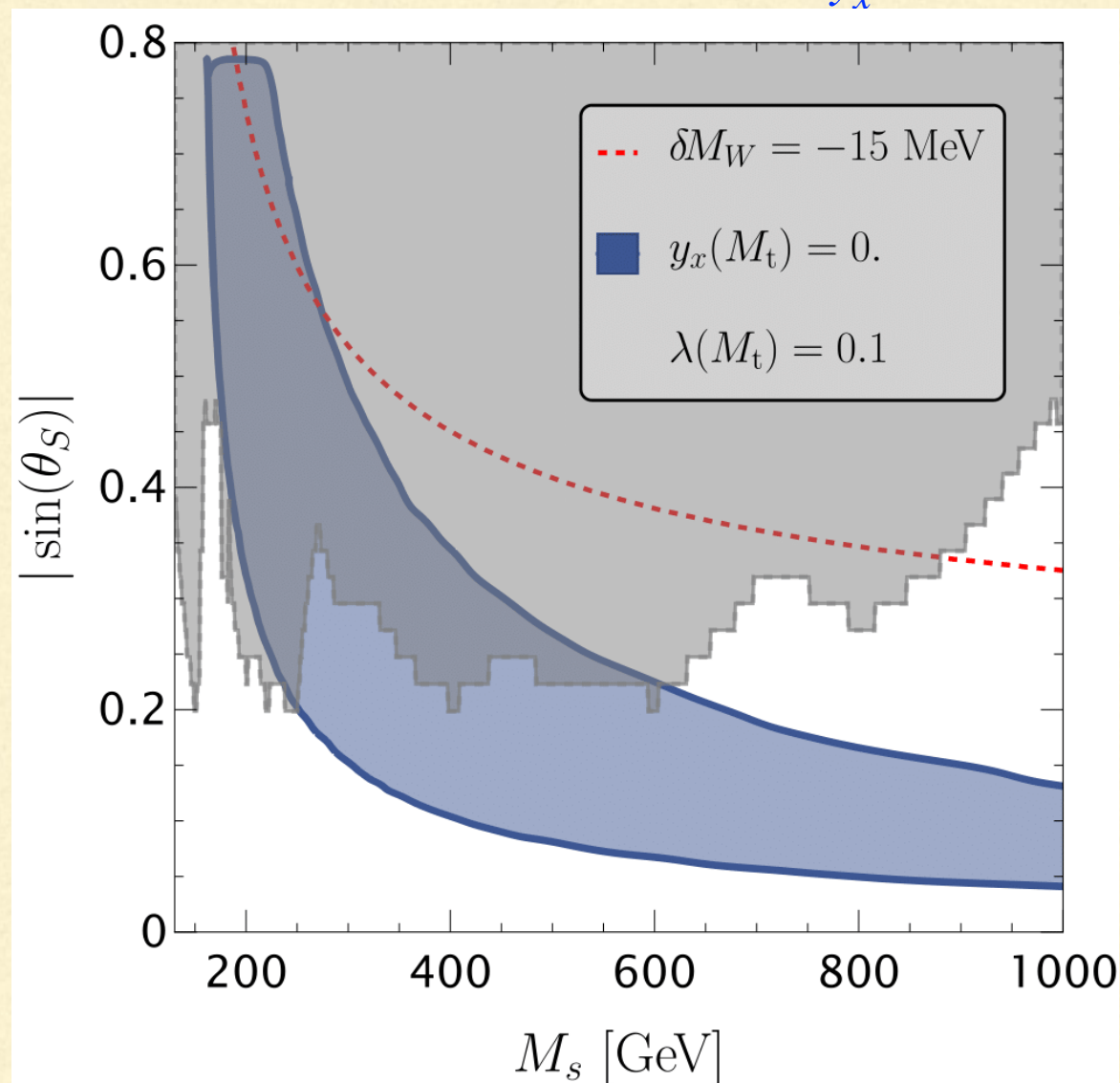
the observation of the Z' or S in the allowed region

Experimental constraints in the scalar sector from direct searches and M_W

■ $M_s > M_h$:

[Zoltán Péli and ZT, arXiv: [2204.07100](#)]

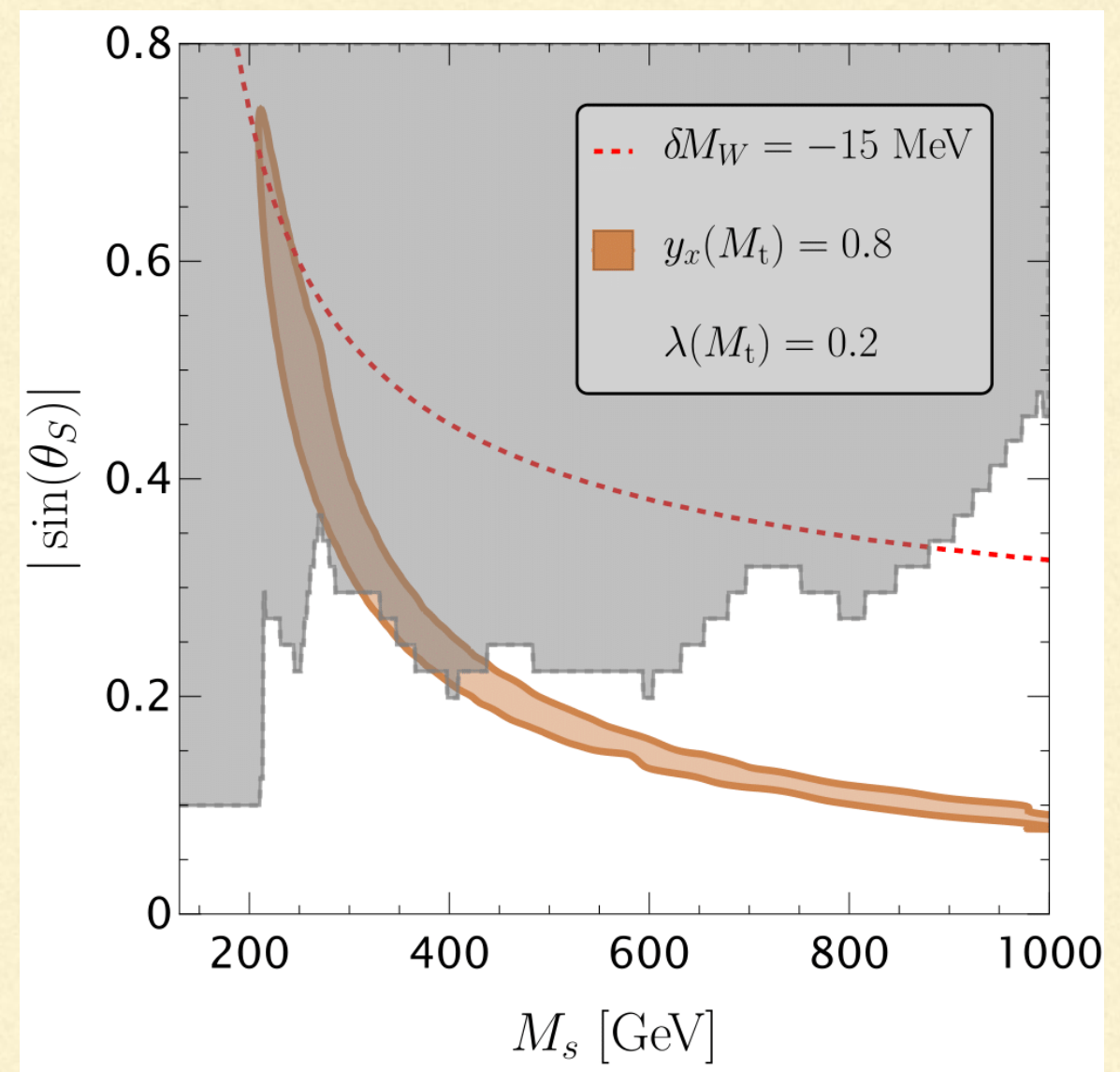
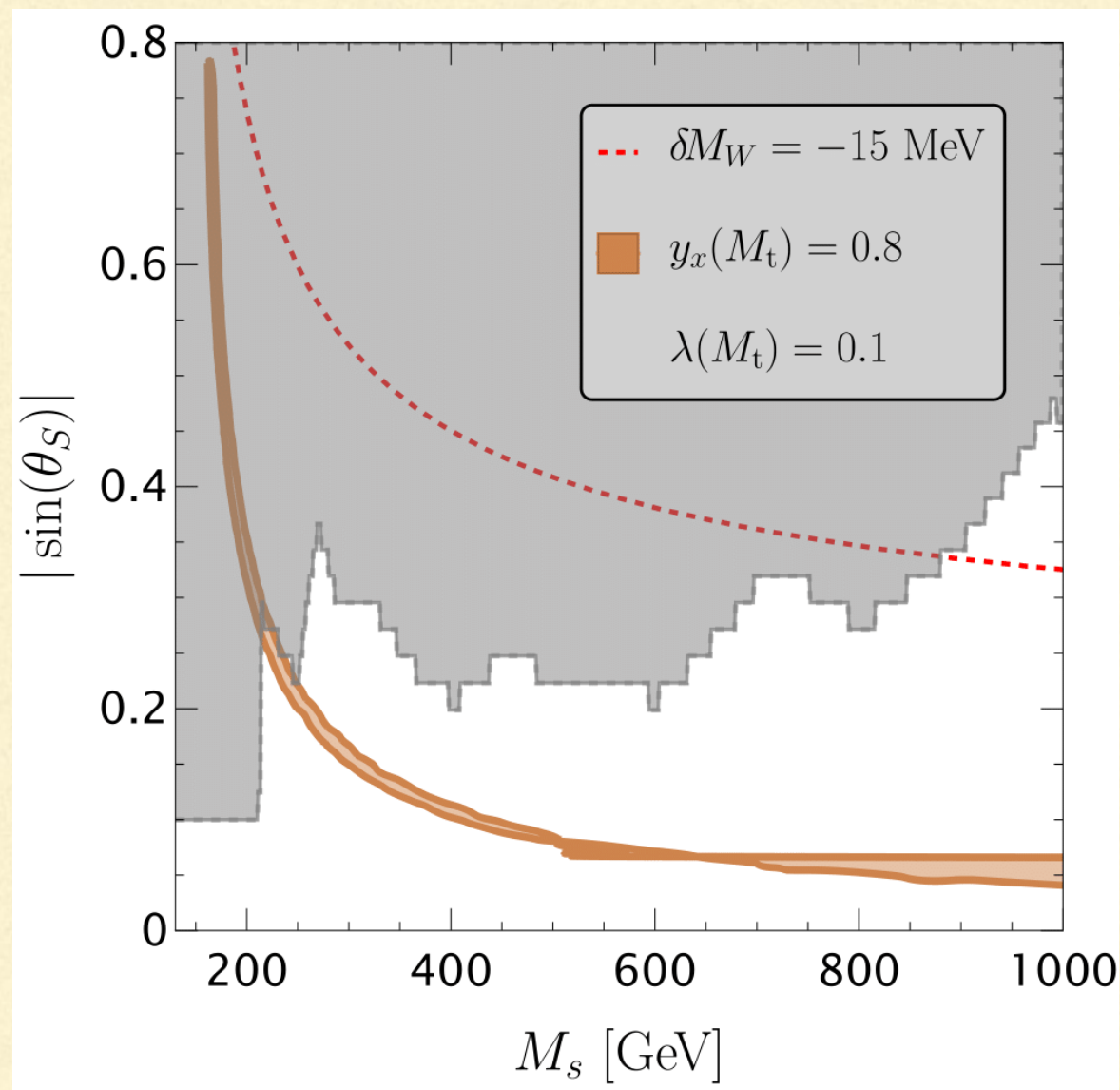
$y_x = 0$: scalar sector decouples



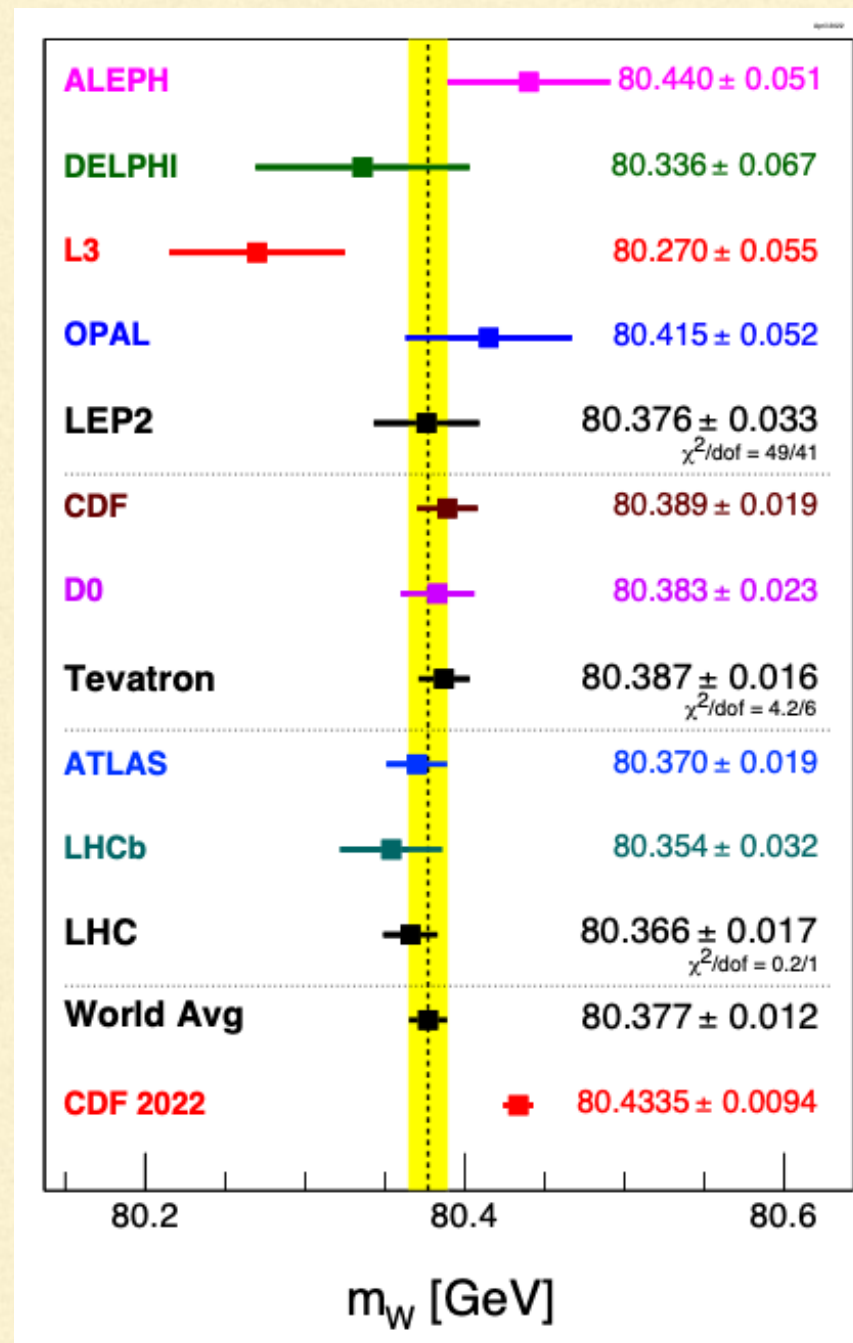
Experimental constraints in the scalar sector from direct searches and M_W

■ $M_s > M_h$:

[Zoltán Péli and ZT, arXiv: [2204.07100](#)]



M_W is measured and computed precisely
(with per myriad precision)



[PDG 2023]

Prediction of M_W in the SWSM

- Can be determined from the decay width of the muon:

$$M_W^2 = \frac{\cos^2 \theta_Z M_Z^2 + \sin^2 \theta_Z M_{Z'}^2}{2} \left(1 + \sqrt{1 - \frac{4\pi\alpha / (\sqrt{2}G_F)}{\cos^2 \theta_Z M_Z^2 + \sin^2 \theta_Z M_{Z'}^2} \frac{1}{1 - \Delta r_{SM} - (\Delta r_{BSM}^{(1)} + \Delta r_{BSM}^{(2)})}} \right)$$

- Valid in $\overline{\text{MS}}$
- θ_Z is the $Z - Z'$ mixing angle
- Δr_{SM} collects the SM quantum corrections (known completely at two loops and partially at three loops)
- $\Delta r_{BSM}^{(1)}$ collects the formally SM quantum corrections but with BSM loops
- $\Delta r_{BSM}^{(2)}$ collects the BSM corrections to $M_{Z'}$ and θ_Z

[Zoltán Péli and ZT, arXiv: [2305.11931](https://arxiv.org/abs/2305.11931)]

Conclusions

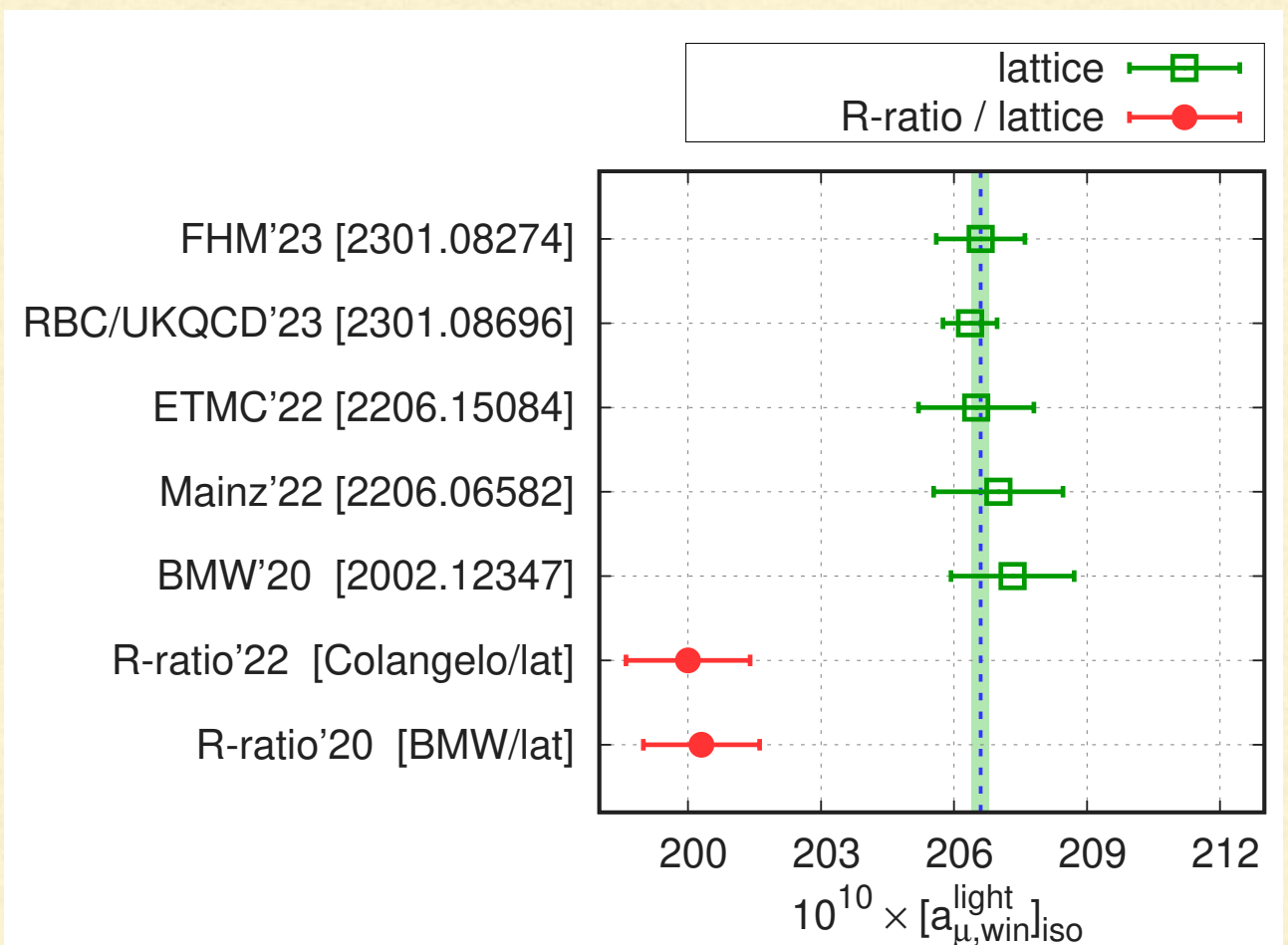
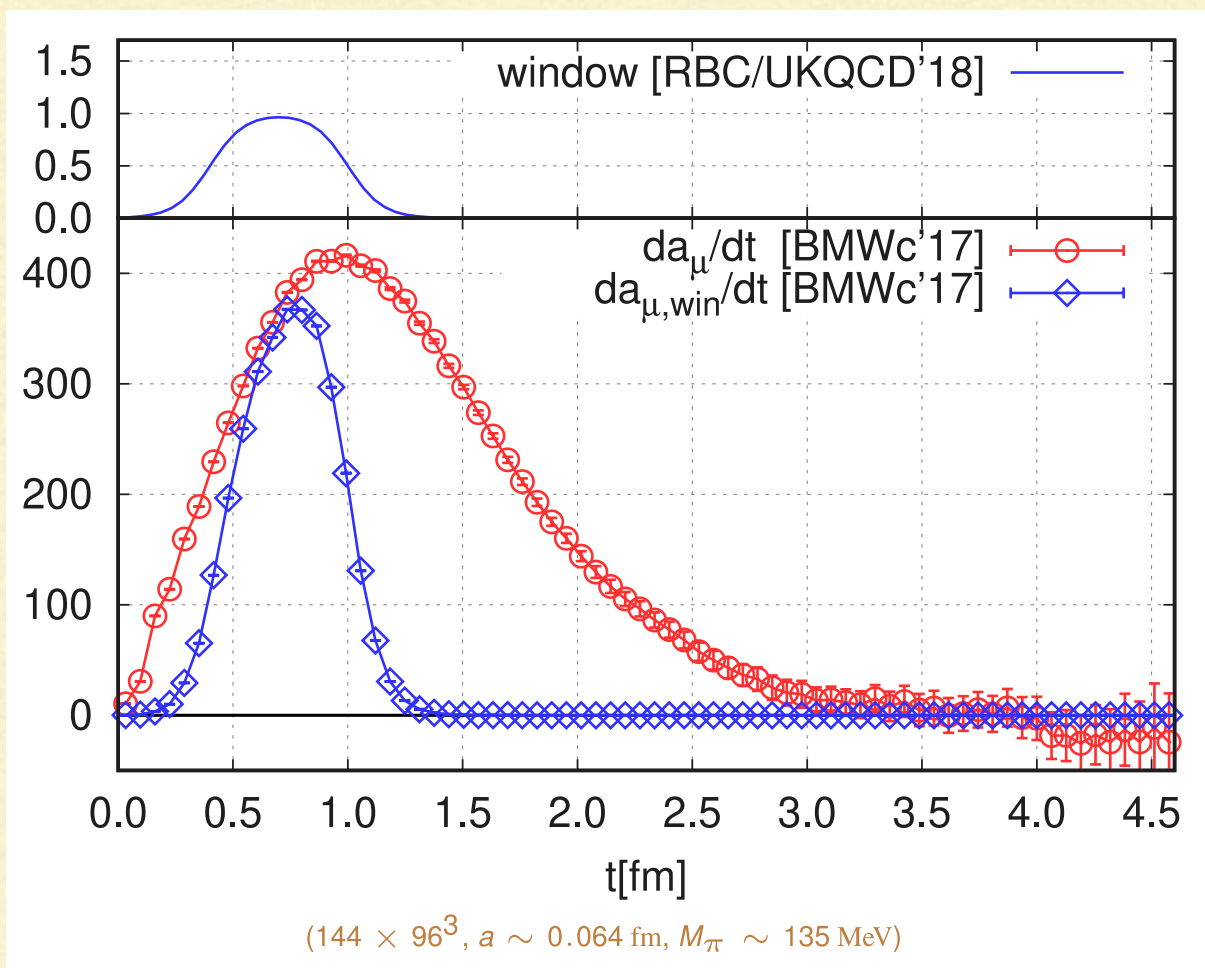
- Established observations require physics beyond SM, but do not suggest rich BSM physics
- $U(1)_Z$ superweak extension has the potential of explaining all known results beyond the SM
- Neutrino masses are generated by SSB at tree level
- One-loop corrections to the tree-level neutrino mass matrix computed and found to be small (below 1‰) in the parameter space relevant in the SWSM
- Lightest sterile neutrino is a candidate DM particle in the [10,50] MeV mass range for freeze-out mechanism with resonant enhancement → predicts an approximate mass relation between vector boson and lightest sterile neutrino
- In the scalar sector we find non-empty parameter space for $M_s > M_h$
- Contributions to EWPOs (e.g. M_W , lepton g-2) are negligible in the superweak region and a systematic exploration of the parameter space is ongoing

the end

Appendix

Status of the muon anomalous magnetic moment: window observable

- restrict correlation window to $[0.4, 1.0]$ fm:
 - two orders of magnitude easier (less CPU, less manpower needed)
- lattice vs. R-ratio: 4.9σ tension:



Non-standard interactions and the SWSM

[Timo J. Kärkäinen and ZT, arXiv: [2301.06621](#)]

$$\mathcal{O}_{6a} = \frac{C_{6a}}{\Lambda^2} (\bar{L} \gamma^\mu P_L L) (\bar{f} \gamma_\mu P_X f)$$

where Λ is the scale of new physics, can be as low as few MeV,
which can be probed in

Coherent Elastic Neutrino-Nucleus Scattering (CEvNS)

Standard parametrization of NSI:

$$\mathcal{L}_{\text{NSI}} = -2\sqrt{2}G_F \sum_{f,X=\pm,\ell,\ell'} \varepsilon_{\ell,\ell'}^{f,X} (\bar{\nu}_\ell \gamma^\mu P_L \nu_{\ell'}) (\bar{f} \gamma_\mu P_X f)$$

where $\varepsilon_{\ell,\ell'}^{f,X} \propto +\frac{1}{q^2}$ if $q^2 \gg M^2$, "light NSI" for a mediator
 $\varepsilon_{\ell,\ell'}^{f,X} \propto -\frac{1}{M^2}$ if $q^2 \ll M^2$, "heavy NSI", of mass M

Non-standard interactions and the SWSM

assume $M = 50 \text{ MeV}$, which is

- light in CHARM or NuTeV $q^2 = O((20 \text{ GeV})^2)$
- heavy in neutrino oscillation experiments $q^2 \approx 0$
- but $q^2 \approx M^2$ in CE ν NS

We can still apply the NSI formalism using the full propagator with q^2 being the characteristic momentum transfer squared

- Can be used to [Timo J. Kärkäinen and ZT, arXiv: [2301.06621](#)]
 - Constrain the parameter space of SWSM
 - Predict relations between NSI couplings assuming SWSM

Non-standard interactions and the SWSM

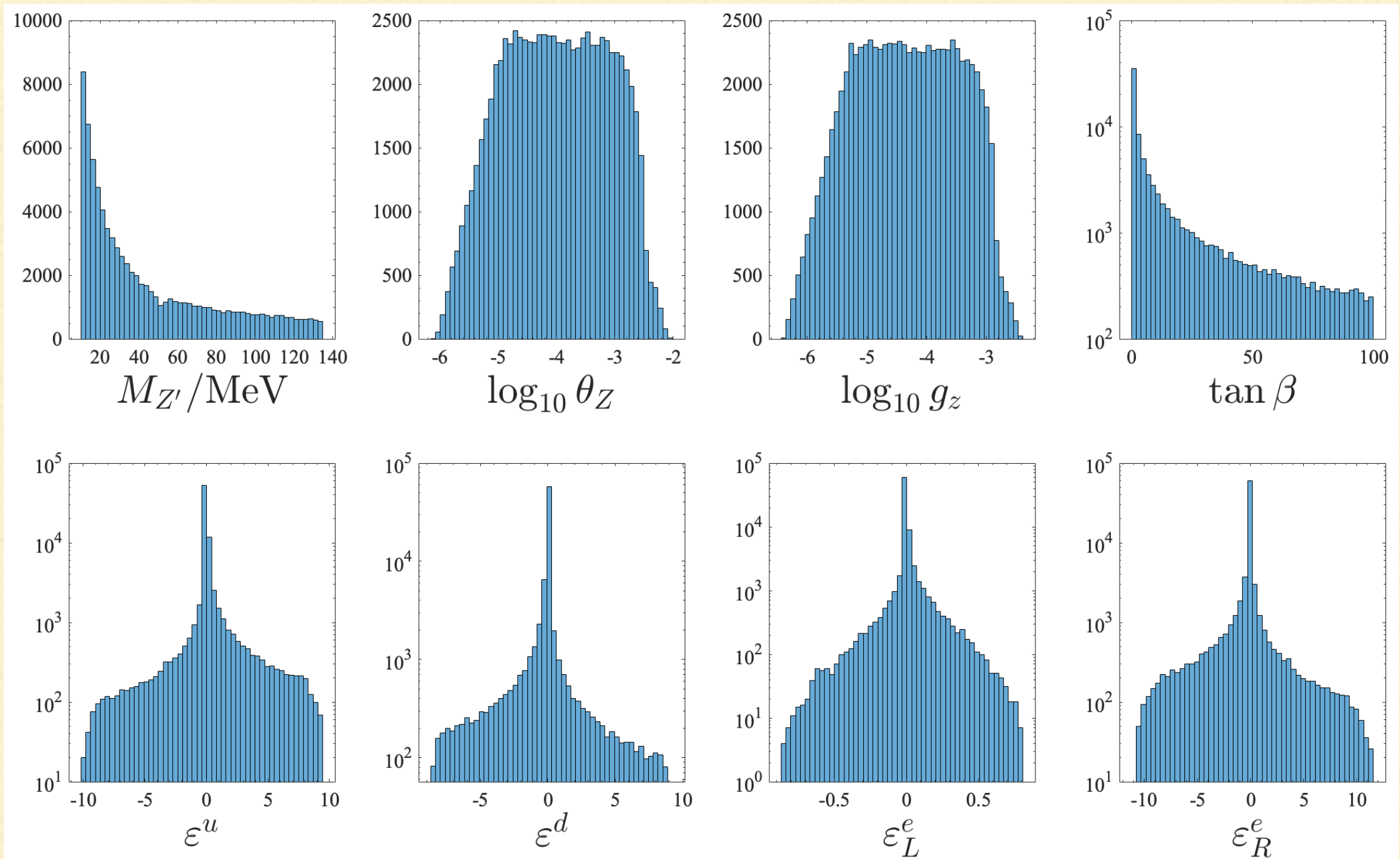
- High-energy theory enforces texture for NSI matrix:

SWSM

$\varepsilon_{\ell\ell}^m = \underbrace{\varepsilon_{\ell\ell}^e + 2\varepsilon_{\ell\ell}^u + \varepsilon_{\ell\ell}^d}_{=0} + \frac{N_n}{N_e}(\varepsilon_{\ell\ell}^u + 2\varepsilon_{\ell\ell}^d)$	$\begin{pmatrix} \varepsilon_{ee}^m & \varepsilon_{e\mu}^m & \varepsilon_{e\tau}^m \\ \varepsilon_{e\mu}^{m*} & \varepsilon_{\mu\mu}^m & \varepsilon_{\mu\tau}^m \\ \varepsilon_{e\tau}^{m*} & \varepsilon_{\mu\tau}^{m*} & \varepsilon_{\tau\tau}^m \end{pmatrix}$	$\begin{pmatrix} \varepsilon_e & 0 & 0 \\ 0 & \varepsilon_\mu & 0 \\ 0 & 0 & \varepsilon_\tau \end{pmatrix}$	$\begin{pmatrix} \varepsilon & 0 & 0 \\ 0 & \varepsilon & 0 \\ 0 & 0 & \varepsilon \end{pmatrix}$
	$\mu - \tau$ symmetry	Flavour-conserving	Flavour-universal
CLFV decays	✓	No	No
ν oscillation	✓	✓	No
CE ν NS	✓	✓	✓
ν scattering	maybe	maybe	maybe

- Existing limits on NSI constrain the parameters of the high-energy theory

Non-standard interactions and the SWSM: preferred regions of the parameters



Non-standard interactions and the SWSM: preferred regions of the parameters

